

Problem 2.24

Check the uncertainty principle for the wave function in Equation 2.132. *Hint:* Calculating $\langle p^2 \rangle$ can be tricky, because the derivative of ψ has a step discontinuity at $x = 0$. You may want to use the result in Problem 2.23(b). *Partial answer:* $\langle p^2 \rangle = (m\alpha/\hbar)^2$.

Solution

Equation 2.132 has the information necessary to construct the wave function.

$$\psi(x) = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2}; \quad E = -\frac{m\alpha^2}{2\hbar^2} \quad (2.132)$$

The product solution is

$$\Psi(x, t) = \psi(x) e^{-iEt/\hbar} = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{im\alpha^2 t/(2\hbar^3)}.$$

Calculate the expectation value of x at time t for this bound state.

$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{\infty} \Psi^*(x, t)(x)\Psi(x, t) dx \\ &= \int_{-\infty}^{\infty} \left[\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{-im\alpha^2 t/(2\hbar^3)} \right] (x) \left[\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{im\alpha^2 t/(2\hbar^3)} \right] dx \\ &= \frac{m\alpha}{\hbar^2} \int_{-\infty}^{\infty} x e^{-2m\alpha|x|/\hbar^2} dx \\ &= 0 \end{aligned}$$

Calculate the expectation value of x^2 at time t .

$$\begin{aligned} \langle x^2 \rangle &= \int_{-\infty}^{\infty} \Psi^*(x, t)(x^2)\Psi(x, t) dx \\ &= \int_{-\infty}^{\infty} \left[\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{-im\alpha^2 t/(2\hbar^3)} \right] (x^2) \left[\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{im\alpha^2 t/(2\hbar^3)} \right] dx \\ &= \frac{m\alpha}{\hbar^2} \int_{-\infty}^{\infty} x^2 e^{-2m\alpha|x|/\hbar^2} dx \\ &= \frac{2m\alpha}{\hbar^2} \int_0^{\infty} x^2 e^{-2m\alpha x/\hbar^2} dx \\ &= \frac{2m\alpha}{\hbar^2} \int_0^{\infty} \frac{\partial^2}{\partial \alpha^2} \left(\frac{\hbar^4}{4m^2} e^{-2m\alpha x/\hbar^2} \right) dx \\ &= \frac{\hbar^2 \alpha}{2m} \frac{d^2}{d\alpha^2} \int_0^{\infty} e^{-2m\alpha x/\hbar^2} dx \\ &= \frac{\hbar^2 \alpha}{2m} \frac{d^2}{d\alpha^2} \left(-\frac{\hbar^2}{2m\alpha} e^{-2m\alpha x/\hbar^2} \Big|_0^{\infty} \right) \\ &= \frac{\hbar^2 \alpha}{2m} \frac{d^2}{d\alpha^2} \left(\frac{\hbar^2}{2m\alpha} \right) \\ &= \frac{\hbar^2 \alpha}{2m} \left(\frac{\hbar^2}{m\alpha^3} \right) = \frac{\hbar^4}{2m^2 \alpha^2} \end{aligned}$$

Consequently, the standard deviation in x at time t is

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \frac{\hbar^2}{\sqrt{2m\alpha}}.$$

Now calculate the expectation values of p and p^2 at time t .

$$\begin{aligned} \langle p \rangle &= \int_{-\infty}^{\infty} \Psi^*(x, t) \left(-i\hbar \frac{\partial}{\partial x} \right) \Psi(x, t) dx & \langle p^2 \rangle &= \int_{-\infty}^{\infty} \Psi^*(x, t) \left(-i\hbar \frac{\partial}{\partial x} \right)^2 \Psi(x, t) dx \\ &= -i\hbar \int_{-\infty}^{\infty} [\psi(x)e^{iEt/\hbar}] \frac{\partial}{\partial x} [\psi(x)e^{-iEt/\hbar}] dx & &= -\hbar^2 \int_{-\infty}^{\infty} [\psi(x)e^{iEt/\hbar}] \frac{\partial^2}{\partial x^2} [\psi(x)e^{-iEt/\hbar}] dx \\ &= -i\hbar \int_{-\infty}^{\infty} \psi(x) \frac{d\psi}{dx} dx & &= -\hbar^2 \int_{-\infty}^{\infty} \psi(x) \frac{d^2\psi}{dx^2} dx \\ &= -i\hbar \int_{-\infty}^{\infty} \frac{d}{dx} \left\{ \frac{1}{2} [\psi(x)]^2 \right\} dx & &= -\hbar^2 \left[\underbrace{\psi(x) \frac{d\psi}{dx}}_{=0} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{d\psi}{dx} \frac{d\psi}{dx} dx \right] \\ &= -i\hbar \left\{ \frac{1}{2} [\psi(x)]^2 \right\} \Big|_{-\infty}^{\infty} & &= \hbar^2 \int_{-\infty}^{\infty} \left(\frac{d\psi}{dx} \right)^2 dx \\ &= 0 & &= \hbar^2 \left[\int_{-\infty}^0 \left(\frac{d\psi}{dx} \right)^2 dx + \int_0^{\infty} \left(\frac{d\psi}{dx} \right)^2 dx \right] \end{aligned}$$

Ehrenfest's theorem confirms this.

$$\langle p \rangle = \frac{d\langle x \rangle}{dt} = 0$$

$$\begin{aligned} &= \hbar^2 \left[\int_{-\infty}^0 \left(\frac{d}{dx} \frac{\sqrt{m\alpha}}{\hbar} e^{m\alpha x/\hbar^2} \right)^2 dx \right. \\ &\quad \left. + \int_0^{\infty} \left(\frac{d}{dx} \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha x/\hbar^2} \right)^2 dx \right] \\ &= \hbar^2 \left[\int_{-\infty}^0 \left(\frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{m\alpha x/\hbar^2} \right)^2 dx \right. \\ &\quad \left. + \int_0^{\infty} \left(-\frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha x/\hbar^2} \right)^2 dx \right] \\ &= \frac{m^3\alpha^3}{\hbar^4} \left(\int_{-\infty}^0 e^{2m\alpha x/\hbar^2} dx + \int_0^{\infty} e^{-2m\alpha x/\hbar^2} dx \right) \\ &= \frac{m^3\alpha^3}{\hbar^4} \left[\frac{\hbar^2}{2m\alpha} e^{2m\alpha x/\hbar^2} \Big|_{-\infty}^0 + \left(-\frac{\hbar^2}{2m\alpha} \right) e^{-2m\alpha x/\hbar^2} \Big|_0^{\infty} \right] \\ &= \frac{m^3\alpha^3}{\hbar^4} \left(\frac{\hbar^2}{2m\alpha} + \frac{\hbar^2}{2m\alpha} \right) \\ &= \frac{m^2\alpha^2}{\hbar^2} \end{aligned}$$

Consequently, the standard deviation in p at time t is

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \frac{m\alpha}{\hbar}.$$

The uncertainty product is

$$\sigma_x \sigma_p = \left(\frac{\hbar^2}{\sqrt{2m\alpha}} \right) \left(\frac{m\alpha}{\hbar} \right) = \frac{\hbar}{\sqrt{2}},$$

which is consistent with Heisenberg's uncertainty principle ($\sigma_x \sigma_p \geq \hbar/2$).

What follows is unnecessary but instructive to go over.

$$\langle p \rangle = -i\hbar \int_{-\infty}^{\infty} \psi(x) \frac{d\psi}{dx} dx \quad \langle p^2 \rangle = -\hbar^2 \int_{-\infty}^{\infty} \psi(x) \frac{d^2\psi}{dx^2} dx$$

Let's say we only got up to here, and we want to evaluate these integrals. The issue is that ψ is in terms of the absolute value sign. Write it in terms of the signum (or sign) function, defined to be -1 if $x < 0$ and 1 if $x > 0$, which can be written in terms of the Heaviside function $\theta(x)$. It's known from Problem 2.23 that $d\theta/dx = \delta(x)$.

$$|x| = \begin{cases} -x & \text{if } x < 0 \\ x & \text{if } x > 0 \end{cases} = x \begin{cases} -1 & \text{if } x < 0 \\ 1 & \text{if } x > 0 \end{cases} = x \operatorname{sgn} x = x[\theta(x) - \theta(-x)]$$

Taking the derivative of the absolute value sign gives

$$\begin{aligned} \frac{d}{dx}|x| &= \frac{d}{dx}\{x[\theta(x) - \theta(-x)]\} \\ &= [\theta(x) - \theta(-x)] + x[\theta'(x) - \theta'(-x)(-1)] \\ &= [\theta(x) - \theta(-x)] + x[\delta(x) + \delta(-x)] \\ &= [\theta(x) - \theta(-x)] + x\left[\delta(x) + \frac{1}{|-1|}\delta(x)\right] \\ &= [\theta(x) - \theta(-x)] + \underbrace{2x\delta(x)}_{=0} \\ &= [\theta(x) - \theta(-x)] \\ &= \operatorname{sgn} x. \end{aligned}$$

As a result,

$$\begin{aligned} \psi(x) &= \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \\ \frac{d\psi}{dx} &= \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \left(-\frac{m\alpha}{\hbar^2} \frac{d}{dx}|x| \right) \\ &= -\frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha|x|/\hbar^2} \operatorname{sgn} x \\ \frac{d^2\psi}{dx^2} &= \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \left(-\frac{m\alpha}{\hbar^2} \frac{d}{dx}|x| \right)^2 + \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \left(-\frac{m\alpha}{\hbar^2} \frac{d^2}{dx^2}|x| \right) \\ &= \frac{\sqrt{m^5\alpha^5}}{\hbar^5} e^{-m\alpha|x|/\hbar^2} (\operatorname{sgn} x)^2 - \frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha|x|/\hbar^2} \frac{d}{dx}[\theta(x) - \theta(-x)] \\ &= \frac{\sqrt{m^5\alpha^5}}{\hbar^5} e^{-m\alpha|x|/\hbar^2} - \frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha|x|/\hbar^2} [2\delta(x)]. \end{aligned}$$

Now evaluate $\langle p \rangle$.

$$\begin{aligned}
 \langle p \rangle &= -i\hbar \int_{-\infty}^{\infty} \psi(x) \frac{d\psi}{dx} dx \\
 &= -i\hbar \int_{-\infty}^{\infty} \left(\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \right) \left(-\frac{\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha|x|/\hbar^2} \operatorname{sgn} x \right) dx \\
 &= i \frac{m^2\alpha^2}{\hbar^3} \int_{-\infty}^{\infty} e^{-2m\alpha|x|/\hbar^2} \operatorname{sgn} x dx \\
 &= i \frac{m^2\alpha^2}{\hbar^3} \left[\int_{-\infty}^0 e^{-2m\alpha(-x)/\hbar^2} (-1) dx + \int_0^{\infty} e^{-2m\alpha(x)/\hbar^2} (1) dx \right] \\
 &= i \frac{m^2\alpha^2}{\hbar^3} \left(-\int_{-\infty}^0 e^{2m\alpha x/\hbar^2} dx + \int_0^{\infty} e^{-2m\alpha x/\hbar^2} dx \right) \\
 &= i \frac{m^2\alpha^2}{\hbar^3} \left[-\left(\frac{\hbar^2}{2m\alpha} \right) e^{2m\alpha x/\hbar^2} \Big|_{-\infty}^0 + \left(-\frac{\hbar^2}{2m\alpha} \right) e^{-2m\alpha x/\hbar^2} \Big|_0^{\infty} \right] \\
 &= i \frac{m^2\alpha^2}{\hbar^3} \left(-\frac{\hbar^2}{2m\alpha} + \frac{\hbar^2}{2m\alpha} \right) \\
 &= 0
 \end{aligned}$$

Finally, evaluate $\langle p^2 \rangle$.

$$\begin{aligned}
 \langle p^2 \rangle &= -\hbar^2 \int_{-\infty}^{\infty} \psi(x) \frac{d^2\psi}{dx^2} dx \\
 &= -\hbar^2 \int_{-\infty}^{\infty} \left(\frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} \right) \left[\frac{\sqrt{m^5\alpha^5}}{\hbar^5} e^{-m\alpha|x|/\hbar^2} - \frac{2\sqrt{m^3\alpha^3}}{\hbar^3} e^{-m\alpha|x|/\hbar^2} \delta(x) \right] dx \\
 &= -\hbar^2 \int_{-\infty}^{\infty} \left[\frac{m^3\alpha^3}{\hbar^6} e^{-2m\alpha|x|/\hbar^2} - \frac{2m^2\alpha^2}{\hbar^4} e^{-2m\alpha|x|/\hbar^2} \delta(x) \right] dx \\
 &= -\hbar^2 \left[\frac{m^3\alpha^3}{\hbar^6} \int_{-\infty}^{\infty} e^{-2m\alpha|x|/\hbar^2} dx - \frac{2m^2\alpha^2}{\hbar^4} \int_{-\infty}^{\infty} e^{-2m\alpha|x|/\hbar^2} \delta(x) dx \right] \\
 &= -\hbar^2 \left[\frac{m^3\alpha^3}{\hbar^6} \left(\int_{-\infty}^0 e^{2m\alpha x/\hbar^2} dx + \int_0^{\infty} e^{-2m\alpha x/\hbar^2} dx \right) - \frac{2m^2\alpha^2}{\hbar^4} e^0 \right] \\
 &= -\hbar^2 \left\{ \frac{m^3\alpha^3}{\hbar^6} \left[\frac{\hbar^2}{2m\alpha} e^{2m\alpha x/\hbar^2} \Big|_{-\infty}^0 + \left(-\frac{\hbar^2}{2m\alpha} \right) e^{-2m\alpha x/\hbar^2} \Big|_0^{\infty} \right] - \frac{2m^2\alpha^2}{\hbar^4} \right\} \\
 &= -\hbar^2 \left[\frac{m^3\alpha^3}{\hbar^6} \left(\frac{\hbar^2}{2m\alpha} + \frac{\hbar^2}{2m\alpha} \right) - \frac{2m^2\alpha^2}{\hbar^4} \right] \\
 &= -\frac{m^3\alpha^3}{\hbar^4} \left(\frac{\hbar^2}{m\alpha} \right) + \frac{2m^2\alpha^2}{\hbar^2} \\
 &= -\frac{m^2\alpha^2}{\hbar^2} + \frac{2m^2\alpha^2}{\hbar^2} \\
 &= \frac{m^2\alpha^2}{\hbar^2}
 \end{aligned}$$