

## Exercise 22

Solve the Neumann problem in the half-plane:  $\Delta u = 0$  in  $\{y > 0\}$ ,  $\partial u / \partial y = h(x)$  on  $\{y = 0\}$  with  $u(x, y)$  bounded at infinity. (*Hint:* Consider the problem satisfied by  $v = \partial u / \partial y$ .)

### Solution

#### The Method of Images

The Poisson equation will be solved in the upper half-plane  $D$  with a Neumann boundary condition.

$$\begin{aligned}\Delta U &= f(x, y), & -\infty < x < \infty, & y > 0 \\ \frac{\partial U}{\partial y}(x, 0) &= h(x)\end{aligned}$$

A Neumann function representation for the solution can be obtained from Green's second identity,

$$\iint_D (u \Delta v - v \Delta u) dA = \int_{\text{bdy } D} \left( u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) ds.$$

Let  $u = U(x, y)$  and let  $v = N = N(x, y; x_0, y_0)$  be the Neumann function.

$$\iint_D (U \Delta N - N \Delta U) dA = \int_{\text{bdy } D} \left( U \frac{\partial N}{\partial n} - N \frac{\partial U}{\partial n} \right) ds$$

If we require  $N$  to satisfy

$$\begin{aligned}\Delta N &= \delta(x - x_0) \delta(y - y_0), & -\infty < x < \infty, & y > 0 \\ \frac{\partial N}{\partial n} &= c \quad \text{on bdy } D,\end{aligned}$$

where  $c$  is a constant and  $(x_0, y_0)$  is a point in the upper half-plane, then the identity becomes

$$\iint_D [U(x, y) \delta(x - x_0) \delta(y - y_0) - N(x, y; x_0, y_0) f(x, y)] dA = \int_{\text{bdy } D} \left[ U(x, y) c - N(x, y; x_0, y_0) \frac{\partial U}{\partial n} \right] ds.$$

Write the normal derivative as  $\partial U / \partial n = \nabla U \cdot \hat{\mathbf{n}}$  and split up the integrals.

$$\begin{aligned}\iint_D U(x, y) \delta(x - x_0) \delta(y - y_0) dA - \iint_D N(x, y; x_0, y_0) f(x, y) dA \\ = c \int_{\text{bdy } D} U(x, y) ds - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} ds\end{aligned}$$

The integral involving the delta functions is  $U(x_0, y_0)$ .

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) dA = c \int_{\text{bdy } D} U(x, y) ds - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} ds \quad (1)$$

Determine the constant  $c$  by setting  $u = 1$  and  $v = N(x, y; x_0, y_0)$  in Green's second identity.

$$\begin{aligned}\iint_D \Delta N \, dA &= \int_{\text{bdy } D} \frac{\partial N}{\partial n} \, ds \\ \iint_D \delta(x - x_0)\delta(y - y_0) \, dA &= \int_{\text{bdy } D} c \, ds \\ 1 &= c \int_{\text{bdy } D} ds\end{aligned}$$

Solve for  $c$ .

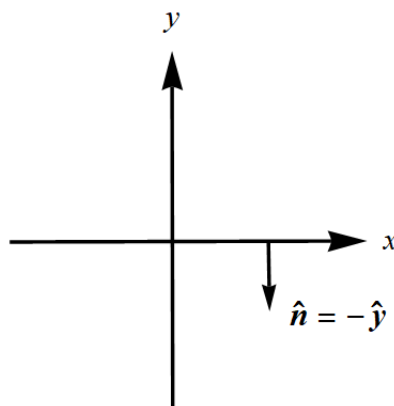
$$c = \frac{1}{\int_{\text{bdy } D} ds} = \frac{1}{\int_{-\infty}^{\infty} dx} = 0$$

Then equation (1) becomes

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA = \frac{\int_{\text{bdy } D} U(x, y) \, ds}{\int_{\text{bdy } D} ds} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} \, ds.$$

The first term on the right side is the average value of  $U$  on the boundary of the upper half-plane ( $y = 0$ ), a constant. Denote it as  $\bar{U}$ .

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA = \bar{U} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} \, ds$$



As the figure illustrates, the outward unit vector normal to the upper half-plane is  $\hat{\mathbf{n}} = -\hat{\mathbf{y}}$ .

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA = \bar{U} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot (-\hat{\mathbf{y}}) \, ds$$

Evaluate the dot product.

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) dA = \bar{U} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \left( -\frac{\partial U}{\partial y} \right) ds$$

Substitute the prescribed boundary condition and write the integration limits.

$$U(x_0, y_0) - \int_0^\infty \int_{-\infty}^\infty N(x, y; x_0, y_0) f(x, y) dx dy = \bar{U} + \int_{-\infty}^\infty N(x, y; x_0, y_0)|_{y=0} h(x) dx$$

Solve for  $U$ .

$$U(x_0, y_0) = \bar{U} + \int_0^\infty \int_{-\infty}^\infty N(x, y; x_0, y_0) f(x, y) dx dy + \int_{-\infty}^\infty N(x, y; x_0, y_0)|_{y=0} h(x) dx$$

Switch the roles of  $x$  and  $y$  with those of  $x_0$  and  $y_0$ , respectively.

$$U(x, y) = \bar{U} + \int_0^\infty \int_{-\infty}^\infty N(x_0, y_0; x, y) f(x_0, y_0) dx_0 dy_0 + \int_{-\infty}^\infty N(x_0, y_0; x, y)|_{y_0=0} h(x_0) dx_0 \quad (2)$$

The Neumann function will now be shown to be symmetric if  $c = 0$ . Set  $u = N(x, y; x_1, y_1)$  and  $v = N(x, y; x_2, y_2)$  in Green's second identity,

$$\begin{aligned} \iint_D [N(x, y; x_1, y_1) \Delta N(x, y; x_2, y_2) - N(x, y; x_2, y_2) \Delta N(x, y; x_1, y_1)] dA \\ = \int_{\text{bdy } D} \left[ N(x, y; x_1, y_1) \frac{\partial N}{\partial n}(x, y; x_2, y_2) - N(x, y; x_2, y_2) \frac{\partial N}{\partial n}(x, y; x_1, y_1) \right] ds, \end{aligned}$$

where  $(x_1, y_1)$  and  $(x_2, y_2)$  are points in  $D$ , and  $N(x, y; x_1, y_1)$  and  $N(x, y; x_2, y_2)$  satisfy

$$\begin{aligned} \Delta N &= \delta(x - x_1) \delta(y - y_1) & \text{in } D & & \Delta N &= \delta(x - x_2) \delta(y - y_2) & \text{in } D \\ \frac{\partial N}{\partial n}(x, y; x_1, y_1) &= c & \text{on bdy } D & & \frac{\partial N}{\partial n}(x, y; x_2, y_2) &= c & \text{on bdy } D. \end{aligned}$$

Substitute these results into the identity.

$$\begin{aligned} \iint_D [N(x, y; x_1, y_1) \delta(x - x_2) \delta(y - y_2) - N(x, y; x_2, y_2) \delta(x - x_1) \delta(y - y_1)] dA \\ = \int_{\text{bdy } D} [N(x, y; x_1, y_1) c - N(x, y; x_2, y_2) c] ds \end{aligned}$$

Split up the integrals on the left and bring  $c$  in front of the integral on the right.

$$\begin{aligned} \iint_D N(x, y; x_1, y_1) \delta(x - x_2) \delta(y - y_2) dA - \iint_D N(x, y; x_2, y_2) \delta(x - x_1) \delta(y - y_1) dA \\ = c \int_{\text{bdy } D} [N(x, y; x_1, y_1) - N(x, y; x_2, y_2)] ds \end{aligned}$$

If  $c = 0$ , then the right side is zero. Evaluate the integrals on the left.

$$N(x_2, y_2; x_1, y_1) - N(x_1, y_1; x_2, y_2) = 0$$

Therefore,  $N(x_2, y_2; x_1, y_1) = N(x_1, y_1; x_2, y_2)$ , and equation (2) becomes

$$U(x, y) = \bar{U} + \int_0^\infty \int_{-\infty}^\infty N(x, y; x_0, y_0) f(x_0, y_0) dx_0 dy_0 + \int_{-\infty}^\infty N(x, y; x_0, y_0)|_{y_0=0} h(x_0) dx_0.$$

The solution for Poisson's equation is known, then, if the Neumann function in the upper half-plane can be determined. Begin by finding the Neumann function in the whole plane (no boundaries).

$$\Delta \mathcal{N} = \delta(x - x_0)\delta(y - y_0), \quad -\infty < x < \infty, \quad -\infty < y < \infty$$

$\mathcal{N}$  can be interpreted as the electrostatic potential, and  $\delta(x - x_0)\delta(y - y_0)$  can be interpreted as the charge density for a unit charge located at  $(x_0, y_0)$ . Since there are no boundaries,  $\mathcal{N}$  is expected to vary solely as a function of the radial distance from  $(x_0, y_0)$ :  $\mathcal{N} = \mathcal{N}(\boldsymbol{z})$ , where  $\boldsymbol{z} = \sqrt{(x - x_0)^2 + (y - y_0)^2}$ . Integrate both sides over a disk centered at  $(x_0, y_0)$  with radius  $\boldsymbol{z}$ .

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \Delta \mathcal{N} dA = \iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \delta(x - x_0)\delta(y - y_0) dA$$

Since the disk contains  $(x_0, y_0)$ , the right side is 1. Write the Laplacian operator as  $\Delta = \nabla^2$

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \nabla^2 \mathcal{N} dA = 1$$

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \nabla \cdot \nabla \mathcal{N} dA = 1$$

and apply the two-dimensional divergence theorem.

$$\int_{(x-x_0)^2+(y-y_0)^2 = \boldsymbol{z}^2} \nabla \mathcal{N} \cdot \hat{\boldsymbol{z}} ds = 1$$

Here  $\hat{\boldsymbol{z}}$  is the unit vector normal to this disk at every point on the boundary.

$$\int_{(x-x_0)^2+(y-y_0)^2 = \boldsymbol{z}^2} \frac{d\mathcal{N}}{d\boldsymbol{z}} ds = 1$$

Because  $\mathcal{N}$  only depends on  $\boldsymbol{z}$ , its derivative is constant on the disk's boundary.

$$\frac{d\mathcal{N}}{d\boldsymbol{z}} \int_{(x-x_0)^2+(y-y_0)^2 = \boldsymbol{z}^2} ds = 1$$

This line integral is just the disk's circumference.

$$\frac{d\mathcal{N}}{d\boldsymbol{z}} (2\pi\boldsymbol{z}) = 1$$

Divide both sides by  $2\pi z$ .

$$\frac{d\mathcal{N}}{dz} = \frac{1}{2\pi z}$$

Integrate both sides with respect to  $z$ .

$$\mathcal{N}(z) = \frac{1}{2\pi} \ln z$$

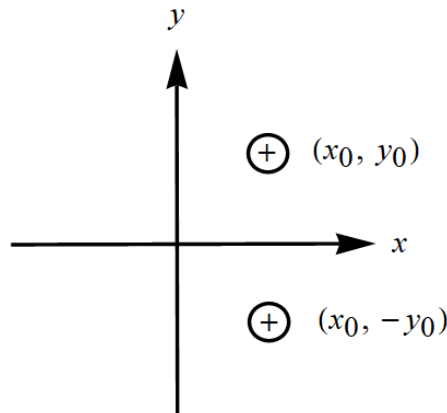
The Neumann function for the whole plane is then

$$\mathcal{N}(x, y; x_0, y_0) = \frac{1}{2\pi} \ln \sqrt{(x - x_0)^2 + (y - y_0)^2}.$$

Now that it's known, the Neumann function for the upper half-plane can be determined by the method of images. A convection of point charges in the whole plane will be arranged so that the boundary condition,

$$\frac{\partial N}{\partial n} = c = \frac{1}{\int_{\text{bdy } D} ds} = \frac{1}{\int_{-\infty}^{\infty} dx} = 0 \text{ on bdy } D \Rightarrow \frac{\partial N}{\partial y}(x, 0) = 0,$$

is satisfied. This derivative of potential can be interpreted as the  $y$ -component of the electric field. For a positive unit charge at  $(x_0, y_0)$ , place another positive unit charge at  $(x_0, -y_0)$  so that every point on the  $x$ -axis is equally spaced from both.



The upper half-plane Neumann function can now be written.

$$N(x, y; x_0, y_0) = +\mathcal{N}(x, y; x_0, y_0) + \mathcal{N}(x, y; x_0, -y_0), \quad y > 0$$

Because  $\mathcal{N}$  is defined over the whole plane, it's important to note the restriction to  $y > 0$  for  $N$ .

$$\begin{aligned} N(x, y; x_0, y_0) &= \frac{1}{2\pi} \ln \sqrt{(x - x_0)^2 + (y - y_0)^2} + \frac{1}{2\pi} \ln \sqrt{(x - x_0)^2 + (y + y_0)^2} \\ &= \frac{1}{2\pi} \ln \left[ \sqrt{(x - x_0)^2 + (y - y_0)^2} \sqrt{(x - x_0)^2 + (y + y_0)^2} \right] \\ &= \frac{1}{4\pi} \ln \{ [(x - x_0)^2 + (y - y_0)^2][(x - x_0)^2 + (y + y_0)^2] \} \end{aligned}$$

Now set  $y_0 = 0$  in this result.

$$\begin{aligned} N(x, y; x_0, y_0)|_{y_0=0} &= \frac{1}{4\pi} \ln \{[(x - x_0)^2 + y^2][(x - x_0)^2 + y^2]\} \\ &= \frac{1}{4\pi} \ln[(x - x_0)^2 + y^2]^2 \\ &= \frac{1}{2\pi} \ln[(x - x_0)^2 + y^2] \end{aligned}$$

Therefore, the solution to Poisson's equation in the upper half-plane with a Neumann boundary condition is

$$\begin{aligned} U(x, y) = \bar{U} + \frac{1}{4\pi} \int_0^\infty \int_{-\infty}^\infty f(x_0, y_0) \ln \{[(x - x_0)^2 + (y - y_0)^2][(x - x_0)^2 + (y + y_0)^2]\} dx_0 dy_0 \\ + \frac{1}{2\pi} \int_{-\infty}^\infty h(x_0) \ln[(x - x_0)^2 + y^2] dx_0. \end{aligned}$$

If  $f = 0$ , then the solution reduces to

$$U(x, y) = \bar{U} + \frac{1}{2\pi} \int_{-\infty}^\infty h(x_0) \ln[(x - x_0)^2 + y^2] dx_0.$$

This answer is in disagreement with the one at the back of the book.

### Using Exercise 6

This answer will be verified by solving the Laplace equation in the upper half-plane with a Neumann boundary condition in a different way.

$$\begin{aligned} \Delta U &= 0, \quad -\infty < x < \infty, \quad y > 0 \\ \frac{\partial U}{\partial y}(x, 0) &= h(x) \end{aligned}$$

Differentiate both sides of the PDE with respect to  $y$ .

$$\begin{aligned} \frac{\partial}{\partial y}(\Delta U) &= \frac{\partial}{\partial y}(0) \\ \frac{\partial}{\partial y} \left( \frac{\partial^2 U}{\partial x^2} + \frac{\partial^2 U}{\partial y^2} \right) &= 0 \\ \frac{\partial}{\partial y} \left( \frac{\partial^2 U}{\partial x^2} \right) + \frac{\partial}{\partial y} \left( \frac{\partial^2 U}{\partial y^2} \right) &= 0 \end{aligned}$$

The mixed derivatives are equal by Clairaut's theorem.

$$\frac{\partial^2}{\partial x^2} \left( \frac{\partial U}{\partial y} \right) + \frac{\partial^2}{\partial y^2} \left( \frac{\partial U}{\partial y} \right) = 0$$

Make the substitution  $V(x, y) = \partial U / \partial y$  in the PDE and its associated boundary condition.

$$\begin{aligned} \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} &= 0 \\ V(x, 0) &= h(x) \end{aligned}$$

The problem becomes the same one that was solved in Exercise 6.

$$V(x, y) = \frac{y}{\pi} \int_{-\infty}^{\infty} \frac{h(x_0)}{(x_0 - x)^2 + y^2} dx_0$$

Now that the solution is known, change back to  $U$ .

$$\frac{\partial U}{\partial y} = \frac{y}{\pi} \int_{-\infty}^{\infty} \frac{h(x_0)}{(x_0 - x)^2 + y^2} dx_0$$

Integrate both sides partially with respect to  $y$ .

$$\begin{aligned} U(x, y) &= \int^y \left[ \frac{s}{\pi} \int_{-\infty}^{\infty} \frac{h(x_0)}{(x_0 - x)^2 + s^2} dx_0 \right] ds + F(x) \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} \left[ \int^y \frac{s}{(x_0 - x)^2 + s^2} ds \right] h(x_0) dx_0 + F(x) \end{aligned}$$

Let

$$\begin{aligned} w &= (x_0 - x)^2 + s^2 \\ dw &= 2s ds \quad \rightarrow \quad \frac{dw}{2} = s ds. \end{aligned}$$

Consequently,

$$\begin{aligned} U(x, y) &= \frac{1}{\pi} \int_{-\infty}^{\infty} \left[ \int^{(x_0-x)^2+y^2} \frac{1}{w} \frac{dw}{2} \right] h(x_0) dx_0 + F(x) \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} \left\{ \frac{1}{2} \ln[(x_0 - x)^2 + y^2] \right\} h(x_0) dx_0 + F(x) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} h(x_0) \ln[(x_0 - x)^2 + y^2] dx_0 + F(x). \end{aligned}$$