

## Exercise 23

Solve the Neumann problem in the quarter-plane  $\{x > 0, y > 0\}$ .

### Solution

The Poisson equation will be solved in a quarter-plane  $D$  with Neumann boundary conditions.

$$\begin{aligned}\Delta U &= f(x, y), \quad x > 0, y > 0 \\ \frac{\partial U}{\partial x}(0, y) &= g(y) \\ \frac{\partial U}{\partial y}(x, 0) &= h(x)\end{aligned}$$

A Neumann function representation for the solution can be obtained from Green's second identity,

$$\iint_D (u\Delta v - v\Delta u) dA = \int_{\text{bdy } D} \left( u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) ds.$$

Let  $u = U(x, y)$  and let  $v = N = N(x, y; x_0, y_0)$  be the Neumann function.

$$\iint_D (U\Delta N - N\Delta U) dA = \int_{\text{bdy } D} \left( U \frac{\partial N}{\partial n} - N \frac{\partial U}{\partial n} \right) ds$$

If we require  $N$  to satisfy

$$\begin{aligned}\Delta N &= \delta(x - x_0)\delta(y - y_0), \quad x > 0, y > 0 \\ \frac{\partial N}{\partial n} &= c \quad \text{on bdy } D,\end{aligned}$$

where  $c$  is a constant and  $(x_0, y_0)$  is a point in the upper half-plane, then the identity becomes

$$\iint_D [U(x, y)\delta(x - x_0)\delta(y - y_0) - N(x, y; x_0, y_0)f(x, y)] dA = \int_{\text{bdy } D} \left[ U(x, y)c - N(x, y; x_0, y_0) \frac{\partial U}{\partial n} \right] ds.$$

Write the normal derivative as  $\partial U/\partial n = \nabla U \cdot \hat{\mathbf{n}}$  and split up the integrals.

$$\begin{aligned}\iint_D U(x, y)\delta(x - x_0)\delta(y - y_0) dA - \iint_D N(x, y; x_0, y_0)f(x, y) dA \\ = c \int_{\text{bdy } D} U(x, y) ds - \int_{\text{bdy } D} N(x, y; x_0, y_0)\nabla U \cdot \hat{\mathbf{n}} ds\end{aligned}$$

The integral involving the delta functions is  $U(x_0, y_0)$ .

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0)f(x, y) dA = c \int_{\text{bdy } D} U(x, y) ds - \int_{\text{bdy } D} N(x, y; x_0, y_0)\nabla U \cdot \hat{\mathbf{n}} ds \quad (1)$$

Determine the constant  $c$  by setting  $u = 1$  and  $v = N(x, y; x_0, y_0)$  in Green's second identity.

$$\begin{aligned}\iint_D \Delta N \, dA &= \int_{\text{bdy } D} \frac{\partial N}{\partial n} \, ds \\ \iint_D \delta(x - x_0)\delta(y - y_0) \, dA &= \int_{\text{bdy } D} c \, ds \\ 1 &= c \int_{\text{bdy } D} ds\end{aligned}$$

Solve for  $c$ .

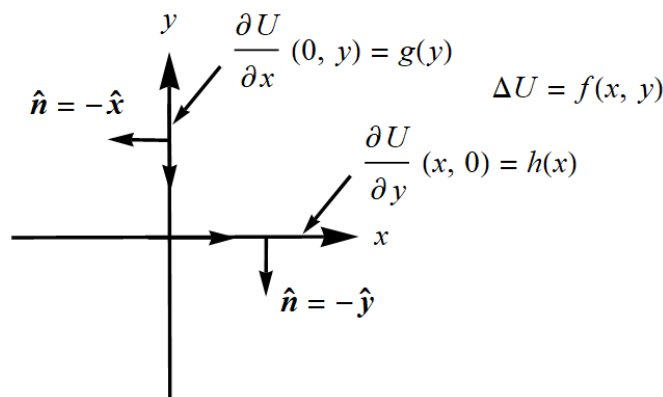
$$c = \frac{1}{\int_{\text{bdy } D} ds} = \frac{1}{\int_{-\infty}^0 (-dy) + \int_0^{\infty} dx} = 0$$

Then equation (1) becomes

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA = \frac{\int_{\text{bdy } D} U(x, y) \, ds}{\int_{\text{bdy } D} ds} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} \, ds.$$

The first term on the right side is the average value of  $U$  on the boundary of the quarter-plane, a constant. Denote it as  $\bar{U}$ .

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA = \bar{U} - \int_{\text{bdy } D} N(x, y; x_0, y_0) \nabla U \cdot \hat{\mathbf{n}} \, ds$$



As the figure illustrates, the outward unit vector normal to the  $y$ -axis is  $\hat{\mathbf{n}} = -\hat{\mathbf{x}}$ , and the outward unit vector normal to the  $x$ -axis is  $\hat{\mathbf{n}} = -\hat{\mathbf{y}}$ .

$$\begin{aligned}U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) \, dA &= \bar{U} - \int_{-\infty}^0 [N(x, y; x_0, y_0) \nabla U \cdot (-\hat{\mathbf{x}})] \Big|_{x=0} (-dy) \\ &\quad - \int_0^{\infty} [N(x, y; x_0, y_0) \nabla U \cdot (-\hat{\mathbf{y}})] \Big|_{y=0} dx\end{aligned}$$

Evaluate the dot products.

$$U(x_0, y_0) - \iint_D N(x, y; x_0, y_0) f(x, y) dA = \bar{U} - \int_0^\infty \left[ N(x, y; x_0, y_0) \left( -\frac{\partial U}{\partial x} \right) \right] \Big|_{x=0} dy \\ - \int_0^\infty \left[ N(x, y; x_0, y_0) \left( -\frac{\partial U}{\partial y} \right) \right] \Big|_{y=0} dx$$

Substitute the prescribed boundary conditions and write the integration limits on the left.

$$U(x_0, y_0) - \int_0^\infty \int_0^\infty N(x, y; x_0, y_0) f(x, y) dx dy = \bar{U} + \int_0^\infty N(x, y; x_0, y_0)|_{x=0} g(y) dy \\ + \int_0^\infty N(x, y; x_0, y_0)|_{y=0} h(x) dx$$

Solve for  $U$ .

$$U(x_0, y_0) = \bar{U} + \int_0^\infty \int_0^\infty N(x, y; x_0, y_0) f(x, y) dx dy + \int_0^\infty N(x, y; x_0, y_0)|_{x=0} g(y) dy \\ + \int_0^\infty N(x, y; x_0, y_0)|_{y=0} h(x) dx$$

Switch the roles of  $x$  and  $y$  with those of  $x_0$  and  $y_0$ , respectively.

$$U(x, y) = \bar{U} + \int_0^\infty \int_0^\infty N(x_0, y_0; x, y) f(x_0, y_0) dx_0 dy_0 + \int_0^\infty N(x_0, y_0; x, y)|_{x_0=0} g(y_0) dy_0 \\ + \int_0^\infty N(x_0, y_0; x, y)|_{y_0=0} h(x_0) dx_0 \quad (2)$$

The Neumann function will now be shown to be symmetric if  $c = 0$ . Set  $u = N(x, y; x_1, y_1)$  and  $v = N(x, y; x_2, y_2)$  in Green's second identity,

$$\iint_D [N(x, y; x_1, y_1) \Delta N(x, y; x_2, y_2) - N(x, y; x_2, y_2) \Delta N(x, y; x_1, y_1)] dA \\ = \int_{\text{bdy } D} \left[ N(x, y; x_1, y_1) \frac{\partial N}{\partial n}(x, y; x_2, y_2) - N(x, y; x_2, y_2) \frac{\partial N}{\partial n}(x, y; x_1, y_1) \right] ds,$$

where  $(x_1, y_1)$  and  $(x_2, y_2)$  are points in  $D$ , and  $N(x, y; x_1, y_1)$  and  $N(x, y; x_2, y_2)$  satisfy

$$\Delta N = \delta(x - x_1) \delta(y - y_1) \quad \text{in } D \qquad \Delta N = \delta(x - x_2) \delta(y - y_2) \quad \text{in } D \\ \frac{\partial N}{\partial n}(x, y; x_1, y_1) = c \quad \text{on bdy } D \qquad \frac{\partial N}{\partial n}(x, y; x_2, y_2) = c \quad \text{on bdy } D.$$

Substitute these results into the identity.

$$\iint_D [N(x, y; x_1, y_1) \delta(x - x_2) \delta(y - y_2) - N(x, y; x_2, y_2) \delta(x - x_1) \delta(y - y_1)] dA \\ = \int_{\text{bdy } D} [N(x, y; x_1, y_1) c - N(x, y; x_2, y_2) c] ds$$

Split up the integrals on the left and bring  $c$  in front of the integral on the right.

$$\begin{aligned} \iint_D N(x, y; x_1, y_1) \delta(x - x_2) \delta(y - y_2) dA - \iint_D N(x, y; x_2, y_2) \delta(x - x_1) \delta(y - y_1) dA \\ = c \int_{\text{bdy } D} [N(x, y; x_1, y_1) - N(x, y; x_2, y_2)] ds \end{aligned}$$

If  $c = 0$ , then the right side is zero. Evaluate the integrals on the left.

$$N(x_2, y_2; x_1, y_1) - N(x_1, y_1; x_2, y_2) = 0$$

Therefore,  $N(x_2, y_2; x_1, y_1) = N(x_1, y_1; x_2, y_2)$ , and equation (2) becomes

$$\begin{aligned} U(x, y) = \bar{U} + \int_0^\infty \int_0^\infty N(x, y; x_0, y_0) f(x_0, y_0) dx_0 dy_0 + \int_0^\infty N(x, y; x_0, y_0)|_{x_0=0} g(y_0) dy_0 \\ + \int_0^\infty N(x, y; x_0, y_0)|_{y_0=0} h(x_0) dx_0. \end{aligned}$$

The solution for Poisson's equation is known, then, if the Neumann function in the quarter-plane can be determined. Begin by finding the Neumann function in the whole plane (no boundaries).

$$\Delta \mathcal{N} = \delta(x - x_0) \delta(y - y_0), \quad -\infty < x < \infty, \quad -\infty < y < \infty$$

$\mathcal{N}$  can be interpreted as the electrostatic potential, and  $\delta(x - x_0) \delta(y - y_0)$  can be interpreted as the charge density for a unit charge located at  $(x_0, y_0)$ . Since there are no boundaries,  $\mathcal{N}$  is expected to vary solely as a function of the radial distance from  $(x_0, y_0)$ :  $\mathcal{N} = \mathcal{N}(\boldsymbol{z})$ , where  $\boldsymbol{z} = \sqrt{(x - x_0)^2 + (y - y_0)^2}$ . Integrate both sides over a disk centered at  $(x_0, y_0)$  with radius  $\boldsymbol{z}$ .

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \Delta \mathcal{N} dA = \iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \delta(x - x_0) \delta(y - y_0) dA$$

Since the disk contains  $(x_0, y_0)$ , the right side is 1. Write the Laplacian operator as  $\Delta = \nabla^2$

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \nabla^2 \mathcal{N} dA = 1$$

$$\iint_{(x-x_0)^2+(y-y_0)^2 \leq \boldsymbol{z}^2} \nabla \cdot \nabla \mathcal{N} dA = 1$$

and apply the two-dimensional divergence theorem.

$$\int_{(x-x_0)^2+(y-y_0)^2 = \boldsymbol{z}^2} \nabla \mathcal{N} \cdot \hat{\boldsymbol{z}} ds = 1$$

Here  $\hat{\boldsymbol{z}}$  is the unit vector normal to this disk at every point on the boundary.

$$\int_{(x-x_0)^2+(y-y_0)^2 = \boldsymbol{z}^2} \frac{d\mathcal{N}}{d\boldsymbol{z}} ds = 1$$

Because  $\mathcal{N}$  only depends on  $z$ , its derivative is constant on the disk's boundary.

$$\frac{d_z \mathcal{N}}{dz} \int_{(x-x_0)^2+(y-y_0)^2=z^2} ds = 1$$

This line integral is just the disk's circumference.

$$\frac{d_z \mathcal{N}}{dz} (2\pi z) = 1$$

Divide both sides by  $2\pi z$ .

$$\frac{d_z \mathcal{N}}{dz} = \frac{1}{2\pi z}$$

Integrate both sides with respect to  $z$ .

$$\mathcal{N}(z) = \frac{1}{2\pi} \ln z$$

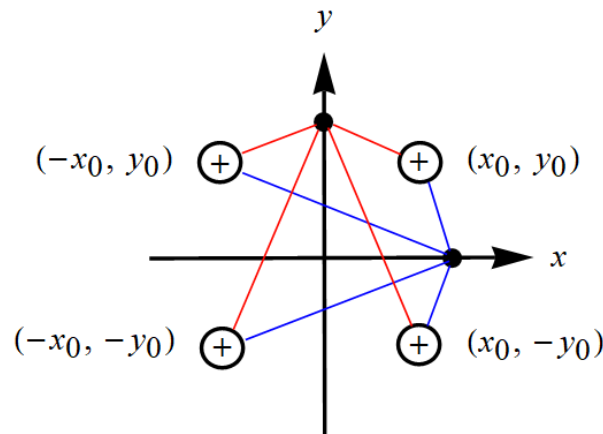
The Neumann function for the whole plane is then

$$\mathcal{N}(x, y; x_0, y_0) = \frac{1}{2\pi} \ln \sqrt{(x-x_0)^2 + (y-y_0)^2}.$$

Now that it's known, the Neumann function for the quarter-plane can be determined by the method of images. A convection of point charges in the whole plane will be arranged so that the boundary condition,

$$\frac{\partial N}{\partial n} = c = \frac{1}{\int_{\text{bdy } D} ds} = \frac{1}{\int_{-\infty}^0 (-dy) + \int_0^{\infty} dx} = 0 \text{ on bdy } D \Rightarrow \begin{cases} \frac{\partial N}{\partial x}(0, y) = 0 \\ \frac{\partial N}{\partial y}(x, 0) = 0 \end{cases},$$

is satisfied. These derivatives of potential with respect to  $x$  and  $y$  can be interpreted as the  $x$ - and  $y$ -components of the electric field, respectively. For a positive unit charge at  $(x_0, y_0)$ , place positive unit charges at  $(x_0, -y_0)$ ,  $(-x_0, y_0)$ , and  $(-x_0, -y_0)$  so that every point on the positive  $x$ - and  $y$ -axes is equally spaced from each pair.



The quarter-plane Neumann function can now be written.

$$N(x, y; x_0, y_0) = +\mathcal{N}(x, y; x_0, y_0) + \mathcal{N}(x, y; x_0, -y_0) + \mathcal{N}(x, y; -x_0, y_0) + \mathcal{N}(x, y; -x_0, -y_0), \quad x > 0, y > 0$$

Because  $\mathcal{N}$  is defined over the whole plane, it's important to note the restriction to  $x > 0, y > 0$  for  $N$ .

$$\begin{aligned} N(x, y; x_0, y_0) &= \frac{1}{2\pi} \ln \sqrt{(x-x_0)^2 + (y-y_0)^2} + \frac{1}{2\pi} \ln \sqrt{(x-x_0)^2 + (y+y_0)^2} + \frac{1}{2\pi} \ln \sqrt{(x+x_0)^2 + (y-y_0)^2} + \frac{1}{2\pi} \ln \sqrt{(x+x_0)^2 + (y+y_0)^2} \\ &= \frac{1}{2\pi} \ln \left[ \sqrt{(x-x_0)^2 + (y-y_0)^2} \sqrt{(x-x_0)^2 + (y+y_0)^2} \sqrt{(x+x_0)^2 + (y-y_0)^2} \sqrt{(x+x_0)^2 + (y+y_0)^2} \right] \\ &= \frac{1}{4\pi} \ln \{ [(x-x_0)^2 + (y-y_0)^2][(x-x_0)^2 + (y+y_0)^2][(x+x_0)^2 + (y-y_0)^2][(x+x_0)^2 + (y+y_0)^2] \} \end{aligned}$$

Now set  $x_0 = 0$  and  $y_0 = 0$  in this result.

$$\begin{aligned} N(x, y; x_0, y_0)|_{x_0=0} &= \frac{1}{4\pi} \ln \{ [x^2 + (y-y_0)^2][x^2 + (y+y_0)^2][x^2 + (y-y_0)^2][x^2 + (y+y_0)^2] \} \\ &= \frac{1}{4\pi} \ln \{ [x^2 + (y-y_0)^2]^2 [x^2 + (y+y_0)^2]^2 \} \\ &= \frac{1}{2\pi} \ln \{ [x^2 + (y-y_0)^2][x^2 + (y+y_0)^2] \} \end{aligned}$$

$$\begin{aligned} N(x, y; x_0, y_0)|_{y_0=0} &= \frac{1}{4\pi} \ln \{ [(x-x_0)^2 + y^2][(x-x_0)^2 + y^2][(x+x_0)^2 + y^2][(x+x_0)^2 + y^2] \} \\ &= \frac{1}{4\pi} \ln \{ [(x-x_0)^2 + y^2]^2 [(x+x_0)^2 + y^2]^2 \} \\ &= \frac{1}{2\pi} \ln \{ [(x-x_0)^2 + y^2][(x+x_0)^2 + y^2] \} \end{aligned}$$

Therefore, the solution to Poisson's equation in the quarter-plane ( $x > 0, y > 0$ ) with Neumann boundary conditions is

$$\begin{aligned} U(x, y) &= \bar{U} + \frac{1}{4\pi} \int_0^\infty \int_0^\infty f(x_0, y_0) \ln \{ [(x-x_0)^2 + (y-y_0)^2][(x-x_0)^2 + (y+y_0)^2][(x+x_0)^2 + (y-y_0)^2][(x+x_0)^2 + (y+y_0)^2] \} dx_0 dy_0 \\ &\quad + \frac{1}{2\pi} \int_0^\infty g(y_0) \ln \{ [x^2 + (y-y_0)^2][x^2 + (y+y_0)^2] \} dy_0 + \frac{1}{2\pi} \int_0^\infty h(x_0) \ln \{ [(x-x_0)^2 + y^2][(x+x_0)^2 + y^2] \} dx_0. \end{aligned}$$