

Problem 3B.7

Momentum fluxes for creeping flow into a slot (Fig. 3B.7). An incompressible Newtonian liquid is flowing very slowly into a thin slot of thickness $2B$ (in the y direction) and width W (in the z direction). The mass rate of flow in the slot is w . From the results of Problem 2B.3 it can be shown that the velocity distribution within the slot is

$$v_x = \frac{3w}{4BW\rho} \left[1 - \left(\frac{y}{B} \right)^2 \right] \quad v_y = 0 \quad v_z = 0 \quad (3B.7-1)$$

at locations not too near the inlet. In the region outside the slot the components of the velocity for *creeping flow* are

$$v_x = -\frac{2w}{\pi W\rho} \frac{x^3}{(x^2 + y^2)^2} \quad (3B.7-2)$$

$$v_y = -\frac{2w}{\pi W\rho} \frac{x^2 y}{(x^2 + y^2)^2} \quad (3B.7-3)$$

$$v_z = 0 \quad (3B.7-4)$$

Equations 3B.7-1 to 4 are only approximate in the region near the slot entry for both $x \geq 0$ and $x \leq 0$.

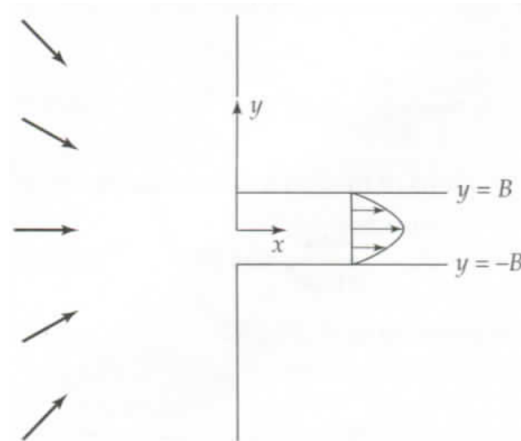


Fig. 3B.7. Flow of a liquid into a slot from a semi-infinite region $x < 0$.

- Find the components of the convective momentum flux $\rho \mathbf{v} \mathbf{v}$ inside and outside the slot.
- Evaluate the xx -component of $\rho \mathbf{v} \mathbf{v}$ at $x = -a$, $y = 0$.
- Evaluate the xy -component of $\rho \mathbf{v} \mathbf{v}$ at $x = -a$, $y = +a$.
- Does the total flow of kinetic energy through the plane $x = -a$ equal the total flow of kinetic energy through the slot?
- Verify that the velocity distributions given in Eqs. 3B.7-1 to 4 satisfy the relation $(\nabla \cdot \mathbf{v}) = 0$.

- (f) Find the normal stress τ_{xx} at the plane $y = 0$ and also on the solid surface at $x = 0$.
- (g) Find the shear stress τ_{yx} on the solid surface at $x = 0$. Is this result surprising? Does sketching the velocity profile v_y vs. x at some plane $y = a$ assist in understanding the result?

Solution

The convective momentum flux, a second-order tensor, is

$$\begin{aligned}\rho\mathbf{v}\mathbf{v} &= \rho \left(\sum_{i=1}^3 v_i \boldsymbol{\delta}_i \right) \left(\sum_{j=1}^3 v_j \boldsymbol{\delta}_j \right) \\ &= \rho(v_1 \boldsymbol{\delta}_1 + v_2 \boldsymbol{\delta}_2 + v_3 \boldsymbol{\delta}_3)(v_1 \boldsymbol{\delta}_1 + v_2 \boldsymbol{\delta}_2 + v_3 \boldsymbol{\delta}_3) \\ &= \rho(v_1^2 \boldsymbol{\delta}_1 \boldsymbol{\delta}_1 + v_1 v_2 \boldsymbol{\delta}_1 \boldsymbol{\delta}_2 + v_1 v_3 \boldsymbol{\delta}_1 \boldsymbol{\delta}_3 \\ &\quad + v_1 v_2 \boldsymbol{\delta}_2 \boldsymbol{\delta}_1 + v_2^2 \boldsymbol{\delta}_2 \boldsymbol{\delta}_2 + v_2 v_3 \boldsymbol{\delta}_2 \boldsymbol{\delta}_3 \\ &\quad + v_1 v_3 \boldsymbol{\delta}_3 \boldsymbol{\delta}_1 + v_2 v_3 \boldsymbol{\delta}_3 \boldsymbol{\delta}_2 + v_3^2 \boldsymbol{\delta}_3 \boldsymbol{\delta}_3).\end{aligned}$$

Part (a)

The components of velocity inside the slot are

$$\begin{aligned}v_x &= \frac{3w}{4BW\rho} \left[1 - \left(\frac{y}{B} \right)^2 \right] \\ v_y &= 0 \\ v_z &= 0,\end{aligned}$$

so the convective momentum flux for $x > 0$ is

$$\begin{aligned}\rho\mathbf{v}\mathbf{v} &= \rho v_1^2 \boldsymbol{\delta}_1 \boldsymbol{\delta}_1 \\ &= \rho \frac{9w^2}{16B^2W^2\rho^2} \left[1 - \left(\frac{y}{B} \right)^2 \right]^2 \boldsymbol{\delta}_x \boldsymbol{\delta}_x \\ &= \frac{9w^2}{16B^2W^2\rho} \left(1 - \frac{y^2}{B^2} \right)^2 \boldsymbol{\delta}_x \boldsymbol{\delta}_x.\end{aligned}$$

On the other hand, the components of velocity outside the slot are

$$\begin{aligned}v_x &= -\frac{2w}{\pi W\rho} \frac{x^3}{(x^2 + y^2)^2} \\ v_y &= -\frac{2w}{\pi W\rho} \frac{x^2 y}{(x^2 + y^2)^2} \\ v_z &= 0,\end{aligned}$$

so the convective momentum flux for $x < 0$ is

$$\begin{aligned}\rho\mathbf{v}\mathbf{v} &= \rho(v_1^2 \boldsymbol{\delta}_1 \boldsymbol{\delta}_1 + v_1 v_2 \boldsymbol{\delta}_1 \boldsymbol{\delta}_2 \\ &\quad + v_1 v_2 \boldsymbol{\delta}_2 \boldsymbol{\delta}_1 + v_2^2 \boldsymbol{\delta}_2 \boldsymbol{\delta}_2).\end{aligned}$$

Substitute the components and expand the right side.

$$\begin{aligned}\rho\mathbf{v}\mathbf{v} &= \rho \left[-\frac{2w}{\pi W\rho} \frac{x^3}{(x^2+y^2)^2} \right]^2 \boldsymbol{\delta}_x\boldsymbol{\delta}_x + \rho \left[-\frac{2w}{\pi W\rho} \frac{x^3}{(x^2+y^2)^2} \right] \left[-\frac{2w}{\pi W\rho} \frac{x^2y}{(x^2+y^2)^2} \right] \boldsymbol{\delta}_x\boldsymbol{\delta}_y \\ &\quad + \rho \left[-\frac{2w}{\pi W\rho} \frac{x^3}{(x^2+y^2)^2} \right] \left[-\frac{2w}{\pi W\rho} \frac{x^2y}{(x^2+y^2)^2} \right] \boldsymbol{\delta}_y\boldsymbol{\delta}_x + \rho \left[-\frac{2w}{\pi W\rho} \frac{x^2y}{(x^2+y^2)^2} \right]^2 \boldsymbol{\delta}_y\boldsymbol{\delta}_y \\ &= \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^6}{(x^2+y^2)^4} \boldsymbol{\delta}_x\boldsymbol{\delta}_x + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^5y}{(x^2+y^2)^4} \boldsymbol{\delta}_x\boldsymbol{\delta}_y \\ &\quad + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^5y}{(x^2+y^2)^4} \boldsymbol{\delta}_y\boldsymbol{\delta}_x + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^4y^2}{(x^2+y^2)^4} \boldsymbol{\delta}_y\boldsymbol{\delta}_y\end{aligned}$$

Therefore,

$$\rho\mathbf{v}\mathbf{v} = \begin{cases} \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^6}{(x^2+y^2)^4} \boldsymbol{\delta}_x\boldsymbol{\delta}_x + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^5y}{(x^2+y^2)^4} \boldsymbol{\delta}_x\boldsymbol{\delta}_y \\ \quad + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^5y}{(x^2+y^2)^4} \boldsymbol{\delta}_y\boldsymbol{\delta}_x + \frac{4w^2}{\pi^2 W^2 \rho} \frac{x^4y^2}{(x^2+y^2)^4} \boldsymbol{\delta}_y\boldsymbol{\delta}_y & \text{if } x < 0. \\ \frac{9w^2}{16B^2W^2\rho} \left(1 - \frac{y^2}{B^2}\right)^2 \boldsymbol{\delta}_x\boldsymbol{\delta}_x & \text{if } x > 0 \end{cases}$$

Note that physically each term represents the convective flux of (2nd delta subscript)-momentum across a plane that is perpendicular to the (1st delta subscript)-direction. Observe that some of them are always positive (even for flows different than the one in this problem). The $\boldsymbol{\delta}_x\boldsymbol{\delta}_x$ fluxes are because the liquid always flows from left to right in the positive x -direction. The $\boldsymbol{\delta}_y\boldsymbol{\delta}_y$ flux is as well because even though the liquid travels up in the positive y -direction only for $y < 0$, the liquid travels down in the negative y -direction for $y > 0$. The $\boldsymbol{\delta}_x\boldsymbol{\delta}_y$ flux is positive when the liquid travels up for $y < 0$ and goes in the positive x -direction but is negative when the liquid travels down for $y > 0$ and goes in the positive x -direction. The $\boldsymbol{\delta}_y\boldsymbol{\delta}_x$ flux is positive when the liquid travels to the right and goes in the positive y -direction ($y < 0$) but is negative when the liquid travels to the right and goes in the negative y -direction ($y > 0$).

Part (b)

$x = -a$ is outside the slot, so the relevant xx -component of $\rho\mathbf{v}\mathbf{v}$ is

$$\frac{4w^2}{\pi^2 W^2 \rho} \frac{x^6}{(x^2+y^2)^4}.$$

Plug in $x = -a$ and $y = 0$ and simplify the result.

$$\begin{aligned}\frac{4w^2}{\pi^2 W^2 \rho} \frac{a^6}{a^8} \\ \frac{4w^2}{\pi^2 a^2 W^2 \rho}\end{aligned}$$

Part (c)

$x = -a$ is outside the slot, so the xy -component of $\rho\mathbf{v}\mathbf{v}$ is

$$\frac{4w^2}{\pi^2 W^2 \rho} \frac{x^5 y}{(x^2 + y^2)^4}$$

Plug in $x = -a$ and $y = a$ and simplify the result.

$$\begin{aligned} & \frac{4w^2}{\pi^2 W^2 \rho} \frac{(-a)^5 (a)}{(2a^2)^4} \\ & \frac{4w^2}{\pi^2 W^2 \rho} \frac{-a^6}{16a^8} \\ & -\frac{w^2}{4\pi^2 a^2 W^2 \rho} \end{aligned}$$

Part (d)

For a liquid in motion, the kinetic energy per unit volume is

$$\frac{1}{2} \rho v^2.$$

Multiply this by the volumetric flow rate (volume per unit time) to get the kinetic energy flow.

$$\frac{1}{2} \rho v^2 \frac{dV}{dt}$$

dV/dt is the dot product of velocity and area.

$$\frac{1}{2} \rho v^2 (\mathbf{v} \cdot \mathbf{A})$$

Actually, because the velocity varies with x and y , it's necessary to integrate over the area that the liquid flows through.

$$\int \frac{1}{2} \rho v^2 (\mathbf{v} \cdot d\mathbf{A})$$

Begin by calculating the kinetic energy flow through the $x = -a$ plane. The unit vector perpendicular to this plane is δ_x , which means the dot product yields the x -component of velocity.

$$\int \frac{1}{2} \rho v^2 v_x dA$$

Integrate over the entire $x = -a$ plane.

$$\int_{-\infty}^{\infty} \frac{1}{2} \rho v^2 v_x (W dy)$$

Bring the constants in front and rewrite v^2 .

$$\frac{\rho W}{2} \int_{-\infty}^{\infty} (v_x^2 + v_y^2 + v_z^2) v_x dy$$

Substitute the components of velocity outside the slot and evaluate the integral.

$$\begin{aligned} & \frac{\rho W}{2} \int_{-\infty}^{\infty} \left[\frac{4w^2}{\pi^2 W^2 \rho^2} \frac{x^6}{(x^2 + y^2)^4} + \frac{4w^2}{\pi^2 W^2 \rho^2} \frac{x^4 y^2}{(x^2 + y^2)^4} + 0 \right] \left[-\frac{2w}{\pi W \rho} \frac{x^3}{(x^2 + y^2)^2} \right] dy \\ & \frac{\rho W}{2} \int_{-\infty}^{\infty} \left[\frac{x^6}{(x^2 + y^2)^4} + \frac{x^4 y^2}{(x^2 + y^2)^4} \right] \left[-\frac{8w^3}{\pi^3 W^3 \rho^3} \frac{x^3}{(x^2 + y^2)^2} \right] dy \\ & - \frac{4w^3}{\pi^3 W^2 \rho^2} \int_{-\infty}^{\infty} \left[\frac{x^9}{(x^2 + y^2)^6} + \frac{x^7 y^2}{(x^2 + y^2)^6} \right] dy \end{aligned}$$

Substitute $x = -a$.

$$\begin{aligned} & -\frac{4w^3}{\pi^3 W^2 \rho^2} \int_{-\infty}^{\infty} \left[\frac{-a^9}{(a^2 + y^2)^6} + \frac{-a^7 y^2}{(a^2 + y^2)^6} \right] dy \\ & \frac{4w^3}{\pi^3 W^2 \rho^2} \left[a^9 \int_{-\infty}^{\infty} \frac{dy}{(a^2 + y^2)^6} + a^7 \int_{-\infty}^{\infty} \frac{y^2}{(a^2 + y^2)^6} dy \right] \\ & \frac{4w^3}{\pi^3 W^2 \rho^2} \left(a^9 \cdot \frac{63\pi}{256a^{11}} + a^7 \cdot \frac{7\pi}{256a^9} \right) \\ & \frac{4w^3}{\pi^3 W^2 \rho^2} \left(\frac{35\pi}{128a^2} \right) \\ & \boxed{\frac{35w^3}{32\pi^2 a^2 W^2 \rho^2}} \end{aligned}$$

Again, this is the flow of kinetic energy through the $x = -a$ plane, which is outside the slot. Now the kinetic energy flow through the inside of the slot will be calculated.

$$\int \frac{1}{2} \rho v^2 (\mathbf{v} \cdot d\mathbf{A})$$

For any plane $x = b > 0$ in the slot, the unit vector perpendicular to it is δ_x ; this makes the dot product yield v_x .

$$\int \frac{1}{2} \rho v^2 v_x dA$$

Integrate over the entire $x = b$ plane.

$$\int_{-B}^B \frac{1}{2} \rho v^2 v_x (W dy)$$

Bring the constants in front and rewrite v^2 .

$$\frac{\rho W}{2} \int_{-B}^B (v_x^2 + v_y^2 + v_z^2) v_x dy$$

Substitute the components of velocity inside the slot.

$$\begin{aligned} & \frac{\rho W}{2} \int_{-B}^B \left[\frac{9w^2}{16B^2 W^2 \rho^2} \left(1 - \frac{y^2}{B^2} \right)^2 + 0 + 0 \right] \left[\frac{3w}{4BW\rho} \left(1 - \frac{y^2}{B^2} \right) \right] dy \\ & \frac{\rho W}{2} \int_{-B}^B \left[\frac{27w^3}{64B^3 W^3 \rho^3} \left(1 - \frac{y^2}{B^2} \right)^3 \right] dy \end{aligned}$$

Bring the constants in front and evaluate the integral.

$$\frac{27w^3}{128B^3W^2\rho^2} \int_{-B}^B \left(1 - \frac{y^2}{B^2}\right)^3 dy$$

$$\frac{27w^3}{128B^3W^2\rho^2} \int_{-B}^B \left(1 - \frac{y^2}{B^2}\right)^3 dy$$

$$\frac{27w^3}{128B^3W^2\rho^2} \left(\frac{32B}{35}\right)$$

$$\boxed{\frac{27w^3}{140B^2W^2\rho^2}}$$

Again, this is the flow of kinetic energy inside the slot. Therefore, the flow of kinetic energy through the plane $x = -a$ is not equal to the flow of kinetic energy through the slot.

Part (e)

Check that $\nabla \cdot \mathbf{v} = 0$ is satisfied both inside and outside the slot. The velocity components are as follows.

Inside The Slot

$$v_x = \frac{3w}{4BW\rho} \left[1 - \left(\frac{y}{B}\right)^2\right]$$

$$v_y = 0$$

$$v_z = 0$$

Outside the slot

$$v_x = -\frac{2w}{\pi W\rho} \frac{x^3}{(x^2 + y^2)^2}$$

$$v_y = -\frac{2w}{\pi W\rho} \frac{x^2y}{(x^2 + y^2)^2}$$

$$v_z = 0$$

Inside the slot, we have

$$\begin{aligned} \nabla \cdot \mathbf{v} &= \left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle \cdot \langle v_x, v_y, v_z \rangle \\ &= \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \\ &= \frac{\partial}{\partial x} \frac{3w}{4BW\rho} \left[1 - \left(\frac{y}{B}\right)^2\right] + \frac{\partial}{\partial y}(0) + \frac{\partial}{\partial z}(0) \\ &= \frac{3w}{4BW\rho} \frac{\partial}{\partial x} \left[1 - \left(\frac{y}{B}\right)^2\right] + 0 + 0 \\ &= \frac{3w}{4BW\rho} (0) \\ &= 0. \end{aligned}$$

Outside the slot, we have

$$\begin{aligned}
 \nabla \cdot \mathbf{v} &= \left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle \cdot \langle v_x, v_y, v_z \rangle \\
 &= \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \\
 &= \frac{\partial}{\partial x} \left[-\frac{2w}{\pi W \rho} \frac{x^3}{(x^2 + y^2)^2} \right] + \frac{\partial}{\partial y} \left[-\frac{2w}{\pi W \rho} \frac{x^2 y}{(x^2 + y^2)^2} \right] + \frac{\partial}{\partial z} (0) \\
 &= -\frac{2w}{\pi W \rho} \frac{\partial}{\partial x} \left[\frac{x^3}{(x^2 + y^2)^2} \right] - \frac{2w}{\pi W \rho} \frac{\partial}{\partial y} \left[\frac{x^2 y}{(x^2 + y^2)^2} \right] + 0 \\
 &= -\frac{2w}{\pi W \rho} \left[\frac{3x^2(x^2 + y^2)^2 - 2(x^2 + y^2) \cdot 2x(x^3)}{(x^2 + y^2)^4} \right] - \frac{2w}{\pi W \rho} \left[\frac{x^2(x^2 + y^2)^2 - 2(x^2 + y^2) \cdot 2y(x^2 y)}{(x^2 + y^2)^4} \right] \\
 &= -\frac{2w}{\pi W \rho} \left[\frac{3x^2(x^2 + y^2) - 4x(x^3)}{(x^2 + y^2)^3} \right] - \frac{2w}{\pi W \rho} \left[\frac{x^2(x^2 + y^2) - 4y(x^2 y)}{(x^2 + y^2)^3} \right] \\
 &= -\frac{2w}{\pi W \rho} \left[\frac{3x^2(x^2 + y^2) - 4x(x^3) + x^2(x^2 + y^2) - 4y(x^2 y)}{(x^2 + y^2)^3} \right] \\
 &= -\frac{2w}{\pi W \rho} \left[\frac{\cancel{3x^4} + \cancel{3x^2 y^2} - \cancel{4x^4} + \cancel{x^4} + \cancel{x^2 y^2} - \cancel{4x^2 y^2}}{(x^2 + y^2)^3} \right] \\
 &= 0.
 \end{aligned}$$

Part (f)

The normal stress τ_{xx} is given on page 843.

$$\tau_{xx} = -\mu \left[2 \frac{\partial v_x}{\partial x} \right] + \left(\frac{2}{3} \mu - \kappa \right) (\nabla \cdot \mathbf{v})$$

It was verified in part (e) that $\nabla \cdot \mathbf{v} = 0$.

$$\tau_{xx} = -2\mu \frac{\partial v_x}{\partial x}$$

Inside the slot v_x is only a function of y , meaning τ_{xx} is zero for $x > 0$. Outside the slot, however, the normal stress is

$$\begin{aligned}
 \tau_{xx} &= -2\mu \frac{\partial}{\partial x} \left[-\frac{2w}{\pi W \rho} \frac{x^3}{(x^2 + y^2)^2} \right] \\
 &= \frac{4\mu w}{\pi W \rho} \frac{\partial}{\partial x} \left[\frac{x^3}{(x^2 + y^2)^2} \right] \\
 &= \frac{4\mu w}{\pi W \rho} \left[\frac{3x^2(x^2 + y^2)^2 - 2(x^2 + y^2) \cdot 2x(x^3)}{(x^2 + y^2)^4} \right] \\
 &= \frac{4\mu w}{\pi W \rho} \left[\frac{3x^2(x^2 + y^2) - 4x(x^3)}{(x^2 + y^2)^3} \right] \\
 &= \frac{4\mu w}{\pi W \rho} \frac{3x^2 y^2 - x^4}{(x^2 + y^2)^3}.
 \end{aligned}$$

Therefore,

$$\tau_{xx} = \begin{cases} \frac{4\mu w}{\pi W \rho} \frac{3x^2 y^2 - x^4}{(x^2 + y^2)^3} & \text{if } x < 0 \\ 0 & \text{if } x > 0 \end{cases}.$$

It follows that the normal stress on the $x = 0$ solid surface is

$$\lim_{x \rightarrow 0} \tau_{xx} = 0,$$

and the normal stress on the $y = 0$ plane is

$$\lim_{y \rightarrow 0} \tau_{xx} = \begin{cases} -\frac{4\mu w}{\pi W \rho x^2} & \text{if } x < 0 \\ 0 & \text{if } x > 0 \end{cases}.$$

Part (g)

The shear stress τ_{yx} is also given on page 843.

$$\tau_{yx} = -\mu \left[\frac{\partial v_y}{\partial x} + \frac{\partial v_x}{\partial y} \right]$$

Inside the slot, we have

$$\begin{aligned} \tau_{yx} &= -\mu \left\{ 0 + \frac{\partial}{\partial y} \left[\frac{3w}{4BW\rho} \left(1 - \frac{y^2}{B^2} \right) \right] \right\} \\ &= -\frac{3\mu w}{4BW\rho} \left(-\frac{2y}{B^2} \right) \\ &= \frac{3\mu w}{2B^3W\rho} y. \end{aligned}$$

Outside the slot, we have

$$\begin{aligned} \tau_{yx} &= -\mu \left\{ \frac{\partial}{\partial x} \left[-\frac{2w}{\pi W \rho} \frac{x^2 y}{(x^2 + y^2)^2} \right] + \frac{\partial}{\partial y} \left[-\frac{2w}{\pi W \rho} \frac{x^3}{(x^2 + y^2)^2} \right] \right\} \\ &= -\mu \left\{ -\frac{2w}{\pi W \rho} \frac{\partial}{\partial x} \left[\frac{x^2 y}{(x^2 + y^2)^2} \right] - \frac{2w}{\pi W \rho} \frac{\partial}{\partial y} \left[\frac{x^3}{(x^2 + y^2)^2} \right] \right\} \\ &= \frac{2\mu w}{\pi W \rho} \left\{ \frac{\partial}{\partial x} \left[\frac{x^2 y}{(x^2 + y^2)^2} \right] + \frac{\partial}{\partial y} \left[\frac{x^3}{(x^2 + y^2)^2} \right] \right\} \\ &= \frac{2\mu w}{\pi W \rho} \left\{ \left[\frac{2xy(x^2 + y^2)^2 - 2(x^2 + y^2) \cdot 2x(x^2 y)}{(x^2 + y^2)^4} \right] + \left[\frac{(0)(x^2 + y^2)^2 - 2(x^2 + y^2) \cdot 2y(x^3)}{(x^2 + y^2)^4} \right] \right\} \\ &= \frac{2\mu w}{\pi W \rho} \left\{ \left[\frac{2xy(x^2 + y^2) - 4x(x^2 y)}{(x^2 + y^2)^3} \right] + \left[\frac{-4y(x^3)}{(x^2 + y^2)^3} \right] \right\} \\ &= \frac{2\mu w}{\pi W \rho} \left[\frac{2xy^3 - 6x^3 y}{(x^2 + y^2)^3} \right] \\ &= \frac{2\mu w}{\pi W \rho} \left[\frac{2xy(y^2 - 3x^2)}{(x^2 + y^2)^3} \right]. \end{aligned}$$

Therefore,

$$\tau_{yx} = \begin{cases} \frac{2\mu w}{\pi W \rho} \left[\frac{2xy(y^2 - 3x^2)}{(x^2 + y^2)^3} \right] & \text{if } x < 0 \\ \frac{3\mu w}{2B^3 W \rho} y & \text{if } x > 0 \end{cases}.$$

It follows that the shear stress on the $x = 0$ solid surface (from the left) is

$$\lim_{x \rightarrow 0^-} \tau_{yx} = 0.$$

In light of the components of velocity outside the slot,

$$\begin{aligned} v_x &= -\frac{2w}{\pi W \rho} \frac{x^3}{(x^2 + y^2)^2} \\ v_y &= -\frac{2w}{\pi W \rho} \frac{x^2 y}{(x^2 + y^2)^2} \\ v_z &= 0, \end{aligned}$$

this result for the shear stress is not surprising because v_y is even in x (meaning that $\partial v_y / \partial x$ is zero at $x = 0$) and v_x has x^3 in the numerator (meaning you get zero when $\partial v_x / \partial y$ is evaluated at $x = 0$). Below is a plot of v_y versus x for $w = W = \rho = 1$ on the $y = 1$ plane to illustrate this first point.

