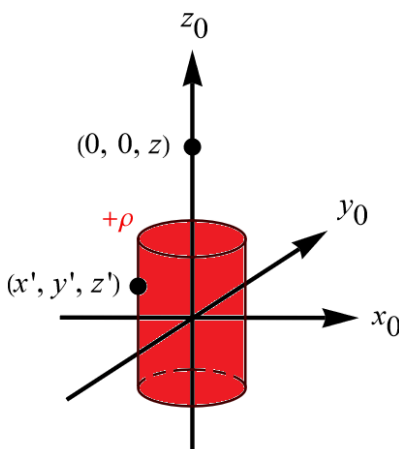


Problem 2.28

Find the potential on the axis of a uniformly charged solid cylinder, a distance z from the center. The length of the cylinder is L , its radius is R , and the charge density is ρ . Use your result to calculate the electric field at this point. (Assume that $z > L/2$.)

Solution

Draw a schematic of the charged cylinder, choosing the origin to be at the center.



$(0, 0, z)$ is where we want to know the electric potential, and (x', y', z') is a point on the charged body. Begin with the basic formula for the electric potential of a continuous volume charge distribution.

$$\begin{aligned} V(\mathbf{r}) &= \frac{1}{4\pi\epsilon_0} \iiint_{\mathcal{V}} \frac{\rho(\mathbf{r}')}{z} d\tau' \\ &= \frac{1}{4\pi\epsilon_0} \iiint_{\mathcal{V}} \frac{\rho}{|\mathbf{r} - \mathbf{r}'|} d\tau' \end{aligned}$$

For the point $(0, 0, z)$ in particular,

$$\begin{aligned} V(0, 0, z) &= \frac{1}{4\pi\epsilon_0} \iiint_{\mathcal{V}} \frac{\rho}{|\langle 0, 0, z \rangle - \langle x', y', z' \rangle|} d\tau' \\ &= \frac{\rho}{4\pi\epsilon_0} \iiint_{\mathcal{V}} \frac{1}{|\langle -x', -y', z - z' \rangle|} d\tau' \\ &= \frac{\rho}{4\pi\epsilon_0} \iiint_{\mathcal{V}} \frac{1}{\sqrt{x'^2 + y'^2 + (z - z')^2}} d\tau'. \end{aligned}$$

Since \mathcal{V} is the cylinder here, switch to cylindrical coordinates.

$$x' = r' \cos \phi'$$

$$y' = r' \sin \phi'$$

$$z' = z'$$

Consequently,

$$\begin{aligned} V(0, 0, z) &= \frac{\rho}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \int_0^{2\pi} \int_0^R \frac{1}{\sqrt{(r' \cos \phi')^2 + (r' \sin \phi')^2 + (z - z')^2}} (r' dr' d\phi' dz') \\ &= \frac{\rho}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \int_0^{2\pi} \left[\int_0^R \frac{r'}{\sqrt{r'^2 + (z - z')^2}} dr' \right] d\phi' dz'. \end{aligned}$$

Make the following substitution.

$$\begin{aligned} u &= r'^2 + (z - z')^2 \\ du &= 2r' dr' \quad \rightarrow \quad \frac{du}{2} = r' dr' \end{aligned}$$

As a result, (note that $z > z'$ because $-L/2 \leq z' \leq L/2$ and $z > L/2$)

$$\begin{aligned} V(0, 0, z) &= \frac{\rho}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \int_0^{2\pi} \left[\int_{0^2 + (z - z')^2}^{R^2 + (z - z')^2} \frac{1}{\sqrt{u}} \left(\frac{du}{2} \right) \right] d\phi' dz' \\ &= \frac{\rho}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \int_0^{2\pi} \left[\sqrt{u} \Big|_{(z - z')^2}^{R^2 + (z - z')^2} \right] d\phi' dz' \\ &= \frac{\rho}{4\pi\epsilon_0} \int_{-L/2}^{L/2} \int_0^{2\pi} \left[\sqrt{R^2 + (z - z')^2} - \sqrt{(z - z')^2} \right] d\phi' dz' \\ &= \frac{\rho}{4\pi\epsilon_0} \left(\int_0^{2\pi} d\phi' \right) \int_{-L/2}^{L/2} \left[\sqrt{R^2 + (z - z')^2} - |z - z'| \right] dz' \\ &= \frac{\rho}{4\pi\epsilon_0} (2\pi) \int_{-L/2}^{L/2} \left[\sqrt{R^2 + (z - z')^2} - (z - z') \right] dz' \\ &= \frac{\rho}{2\epsilon_0} \left[\int_{-L/2}^{L/2} \sqrt{R^2 + (z - z')^2} dz' - z \underbrace{\int_{-L/2}^{L/2} dz'}_{=L} + \underbrace{\int_{-L/2}^{L/2} z' dz'}_{=0} \right]. \end{aligned}$$

Make another substitution.

$$\begin{aligned} z - z' &= R \tan \alpha & \rightarrow & \quad R^2 + (z - z')^2 = R^2 \sec^2 \alpha \\ -dz' &= R \sec^2 \alpha d\alpha & \rightarrow & \quad dz' = -R \sec^2 \alpha d\alpha \end{aligned}$$

So then

$$\begin{aligned} V(0, 0, z) &= \frac{\rho}{2\epsilon_0} \left[-Lz + \int_{\tan^{-1}\left(\frac{z+L/2}{R}\right)}^{\tan^{-1}\left(\frac{z-L/2}{R}\right)} \sqrt{R^2 \sec^2 \alpha} (-R \sec^2 \alpha d\alpha) \right] \\ &= \frac{\rho}{2\epsilon_0} \left[-Lz + R^2 \int_{\tan^{-1}\left(\frac{z-L/2}{R}\right)}^{\tan^{-1}\left(\frac{z+L/2}{R}\right)} \sec^3 \alpha d\alpha \right]. \end{aligned} \tag{1}$$

Use integration by parts to determine the antiderivative of $\sec^3 \alpha$, setting the integration constant to zero.

$$\begin{aligned}
 \int \sec^3 \alpha \, d\alpha &= \int \sec \alpha (\sec^2 \alpha) \, d\alpha \\
 &= \int \sec \alpha \frac{d}{d\alpha} (\tan \alpha) \, d\alpha \\
 &= \sec \alpha \tan \alpha - \int \frac{d}{d\alpha} (\sec \alpha) \tan \alpha \, d\alpha \\
 &= \sec \alpha \tan \alpha - \int (\sec \alpha \tan \alpha) \tan \alpha \, d\alpha \\
 &= \sec \alpha \tan \alpha - \int \sec \alpha \tan^2 \alpha \, d\alpha \\
 &= \sec \alpha \tan \alpha - \int \sec \alpha (\sec^2 \alpha - 1) \, d\alpha \\
 &= \sec \alpha \tan \alpha - \int \sec^3 \alpha \, d\alpha + \int \sec \alpha \, d\alpha \\
 2 \int \sec^3 \alpha \, d\alpha &= \sec \alpha \tan \alpha + \ln |\sec \alpha + \tan \alpha| \\
 \int \sec^3 \alpha \, d\alpha &= \frac{1}{2} (\sec \alpha \tan \alpha + \ln |\sec \alpha + \tan \alpha|)
 \end{aligned}$$

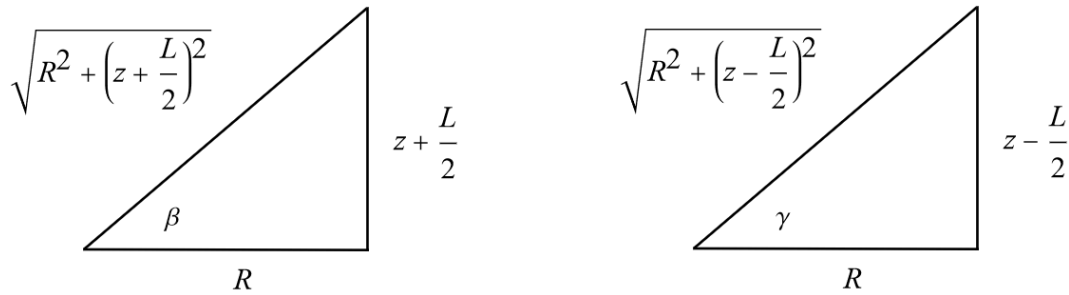
Consequently, equation (1) becomes

$$\begin{aligned}
 V(0, 0, z) &= \frac{\rho}{2\epsilon_0} \left[-Lz + \frac{R^2}{2} (\sec \alpha \tan \alpha + \ln |\sec \alpha + \tan \alpha|) \Big|_{\tan^{-1}\left(\frac{z-L/2}{R}\right)}^{\tan^{-1}\left(\frac{z+L/2}{R}\right)} \right] \\
 &= \frac{\rho}{2\epsilon_0} \left\{ -Lz + \frac{R^2}{2} \left[\sec \tan^{-1} \left(\frac{z+L/2}{R} \right) \tan \tan^{-1} \left(\frac{z+L/2}{R} \right) \right. \right. \\
 &\quad \left. \left. - \sec \tan^{-1} \left(\frac{z-L/2}{R} \right) \tan \tan^{-1} \left(\frac{z-L/2}{R} \right) \right. \right. \\
 &\quad \left. \left. + \ln \left| \frac{\sec \tan^{-1} \left(\frac{z+L/2}{R} \right) + \tan \tan^{-1} \left(\frac{z+L/2}{R} \right)}{\sec \tan^{-1} \left(\frac{z-L/2}{R} \right) + \tan \tan^{-1} \left(\frac{z-L/2}{R} \right)} \right| \right] \right\}. \quad (2)
 \end{aligned}$$

Let

$$\beta = \tan^{-1} \left(\frac{z+L/2}{R} \right) \quad \text{and} \quad \gamma = \tan^{-1} \left(\frac{z-L/2}{R} \right)$$

and draw the implied triangles in order to determine $\sec \beta$, $\tan \beta$, $\sec \gamma$, and $\tan \gamma$.



$$\sec \beta = \frac{\sqrt{R^2 + (z + L/2)^2}}{R}$$

$$\tan \beta = \frac{z + L/2}{R}$$

$$\sec \gamma = \frac{\sqrt{R^2 + (z - L/2)^2}}{R}$$

$$\tan \gamma = \frac{z - L/2}{R}$$

As a result, equation (2) becomes

$$\begin{aligned} V(0, 0, z) &= \frac{\rho}{2\epsilon_0} \left[-Lz + \frac{R^2}{2} \left(\sec \beta \tan \beta - \sec \gamma \tan \gamma + \ln \left| \frac{\sec \beta + \tan \beta}{\sec \gamma + \tan \gamma} \right| \right) \right] \\ &= \frac{\rho}{2\epsilon_0} \left\{ -Lz + \frac{R^2}{2} \left[\frac{\sqrt{R^2 + (z + L/2)^2}}{R} \frac{z + L/2}{R} \right. \right. \\ &\quad \left. \left. - \frac{\sqrt{R^2 + (z - L/2)^2}}{R} \frac{z - L/2}{R} \right. \right. \\ &\quad \left. \left. + \ln \left| \frac{\frac{\sqrt{R^2 + (z + L/2)^2}}{R} + \frac{z + L/2}{R}}{\frac{\sqrt{R^2 + (z - L/2)^2}}{R} + \frac{z - L/2}{R}} \right| \right] \right\} \\ &= \frac{\rho}{2\epsilon_0} \left\{ -Lz + \frac{1}{2} \left[\left(z + \frac{L}{2} \right) \sqrt{R^2 + \left(z + \frac{L}{2} \right)^2} \right. \right. \\ &\quad \left. \left. - \left(z - \frac{L}{2} \right) \sqrt{R^2 + \left(z - \frac{L}{2} \right)^2} \right. \right. \\ &\quad \left. \left. + R^2 \ln \frac{\sqrt{R^2 + \left(z + \frac{L}{2} \right)^2} + \left(z + \frac{L}{2} \right)}{\sqrt{R^2 + \left(z - \frac{L}{2} \right)^2} + \left(z - \frac{L}{2} \right)} \right] \right\}. \end{aligned}$$

Therefore, the electric potential on the axis of a uniformly charged solid cylinder, a distance z from the center, is

$$V(0, 0, z) = \frac{\rho}{4\epsilon_0} \left[-2Lz + \left(z + \frac{L}{2}\right) \sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} - \left(z - \frac{L}{2}\right) \sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} + R^2 \ln \frac{\sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} + \left(z + \frac{L}{2}\right)}{\sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} + \left(z - \frac{L}{2}\right)} \right].$$

Only the z -component of the electric field can be found using

$$\mathbf{E} = -\nabla V$$

because only $V(0, 0, z)$ is known. However, due to the symmetry about the z -axis, the x - and y -components of the electric field are zero.

$$\begin{aligned} \mathbf{E} &= -\frac{d}{dz} V(0, 0, z) \hat{\mathbf{z}} \\ &= \frac{\rho}{2\epsilon_0} \left[L - \sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} + \sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} \right] \hat{\mathbf{z}} \end{aligned}$$