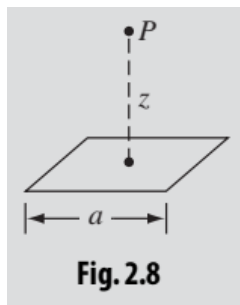


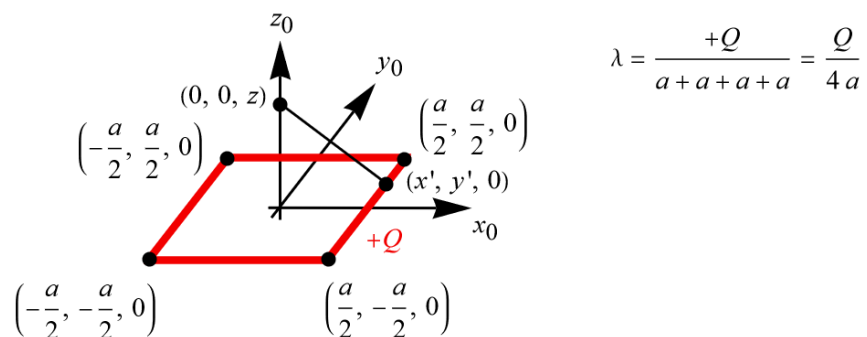
## Problem 2.4

Find the electric field a distance  $z$  above the center of a square loop (side length  $a$ ) carrying a uniform line charge  $\lambda$  (Fig. 2.8). [Hint: Use the result of Ex. 2.2.]



### Solution

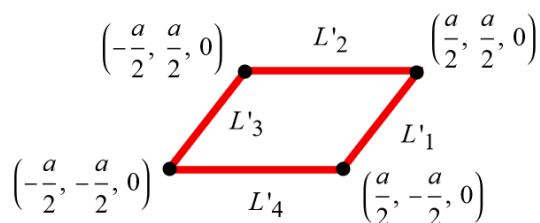
Start by drawing a schematic for some point on the square loop.



The formula for the electric field from a continuous distribution of charge along a line is

$$\begin{aligned} \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\mathbf{r}')}{z^2} \hat{\mathbf{z}} dl' = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^2} \left( \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|} \right) dl' \\ &= \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} (\mathbf{r} - \mathbf{r}') dl', \end{aligned}$$

where the integral is taken over the line where the charge exists. Note that  $\mathbf{r}$  is the position vector to where we want to know the electric field,  $\mathbf{r}'$  is the position vector to the point we chose on the line, and  $z = |\mathbf{r} - \mathbf{r}'|$  is the distance from the point we chose on the line to where we want to know the electric field. Split up the loop into the four straight line segments,  $L'_1$  and  $L'_2$  and  $L'_3$  and  $L'_4$ , shown below.



The parameterizations for these lines are as follows.

$$\begin{aligned} \text{On } L'_1: \quad \mathbf{r}' &= \left\langle \frac{a}{2}, t, 0 \right\rangle, & -\frac{a}{2} \leq t \leq \frac{a}{2} \\ \text{On } L'_2: \quad \mathbf{r}' &= \left\langle t, \frac{a}{2}, 0 \right\rangle, & -\frac{a}{2} \leq t \leq \frac{a}{2} \\ \text{On } L'_3: \quad \mathbf{r}' &= \left\langle -\frac{a}{2}, t, 0 \right\rangle, & -\frac{a}{2} \leq t \leq \frac{a}{2} \\ \text{On } L'_4: \quad \mathbf{r}' &= \left\langle t, -\frac{a}{2}, 0 \right\rangle, & -\frac{a}{2} \leq t \leq \frac{a}{2} \end{aligned}$$

Consequently, the electric field at  $\mathbf{r} = \langle 0, 0, z \rangle$  is

$$\begin{aligned} \mathbf{E} &= \frac{1}{4\pi\epsilon_0} \left[ \int_{L'_1} \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} (\mathbf{r} - \mathbf{r}') dy' + \int_{L'_2} \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} (\mathbf{r} - \mathbf{r}') dx' + \int_{L'_3} \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} (\mathbf{r} - \mathbf{r}') dy' + \int_{L'_4} \frac{\lambda(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} (\mathbf{r} - \mathbf{r}') dx' \right] \\ &= \frac{1}{4\pi\epsilon_0} \left[ \int_{L'_1} \frac{\lambda}{\left[ \sqrt{(0-x')^2 + (0-y')^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle x', y', 0 \rangle) dy' \right. \\ &\quad + \int_{L'_2} \frac{\lambda}{\left[ \sqrt{(0-x')^2 + (0-y')^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle x', y', 0 \rangle) dx' \\ &\quad + \int_{L'_3} \frac{\lambda}{\left[ \sqrt{(0-x')^2 + (0-y')^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle x', y', 0 \rangle) dy' \\ &\quad \left. + \int_{L'_4} \frac{\lambda}{\left[ \sqrt{(0-x')^2 + (0-y')^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle x', y', 0 \rangle) dx' \right] \\ &= \frac{1}{4\pi\epsilon_0} \left\{ \int_{-a/2}^{a/2} \frac{\lambda}{\left[ \sqrt{(0-\frac{a}{2})^2 + (0-t)^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle \frac{a}{2}, t, 0 \rangle) \frac{dy'}{dt} dt \right. \\ &\quad + \int_{-a/2}^{a/2} \frac{\lambda}{\left[ \sqrt{(0-t)^2 + (0-\frac{a}{2})^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle t, \frac{a}{2}, 0 \rangle) \frac{dx'}{dt} dt \\ &\quad + \int_{-a/2}^{a/2} \frac{\lambda}{\left[ \sqrt{(0+\frac{a}{2})^2 + (0-t)^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle -\frac{a}{2}, t, 0 \rangle) \frac{dy'}{dt} dt \\ &\quad \left. + \int_{-a/2}^{a/2} \frac{\lambda}{\left[ \sqrt{(0-t)^2 + (0+\frac{a}{2})^2 + (z-0)^2} \right]^3} (\langle 0, 0, z \rangle - \langle t, -\frac{a}{2}, 0 \rangle) \frac{dx'}{dt} dt \right\}. \end{aligned}$$

Simplify the integrands, combine the integrals, add the vectors, and then integrate the components.

$$\begin{aligned}
 \mathbf{E} &= \frac{\lambda}{4\pi\epsilon_0} \left[ \int_{-a/2}^{a/2} \frac{1}{\left(\frac{a^2}{4} + t^2 + z^2\right)^{3/2}} \left\langle -\frac{a}{2}, -t, z \right\rangle (1) dt + \int_{-a/2}^{a/2} \frac{1}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \left\langle -t, -\frac{a}{2}, z \right\rangle (1) dt \right. \\
 &\quad \left. + \int_{-a/2}^{a/2} \frac{1}{\left(\frac{a^2}{4} + t^2 + z^2\right)^{3/2}} \left\langle \frac{a}{2}, -t, z \right\rangle (1) dt + \int_{-a/2}^{a/2} \frac{1}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \left\langle -t, \frac{a}{2}, z \right\rangle (1) dt \right] \\
 &= \frac{\lambda}{4\pi\epsilon_0} \left[ \int_{-a/2}^{a/2} \frac{1}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \left( \left\langle -\frac{a}{2}, -t, z \right\rangle + \left\langle -t, -\frac{a}{2}, z \right\rangle + \left\langle \frac{a}{2}, -t, z \right\rangle + \left\langle -t, \frac{a}{2}, z \right\rangle \right) dt \right] \\
 &= \frac{\lambda}{4\pi\epsilon_0} \int_{-a/2}^{a/2} \frac{1}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \langle -2t, -2t, 4z \rangle dt \\
 &= \frac{\lambda}{4\pi\epsilon_0} \left\langle \int_{-a/2}^{a/2} \frac{-2t dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}}, \int_{-a/2}^{a/2} \frac{-2t dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}}, \int_{-a/2}^{a/2} \frac{4z dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \right\rangle
 \end{aligned}$$

The first two integrals are zero because the integration interval is symmetric and the integrands are odd functions. The third integral can be made to go from 0 to  $a/2$  by putting a factor of 2 in front because the integrand is an even function.

$$\begin{aligned}
 \mathbf{E} &= \frac{\lambda}{4\pi\epsilon_0} \left\langle 0, 0, 2 \int_0^{a/2} \frac{4z dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \right\rangle \\
 &= \frac{\lambda}{4\pi\epsilon_0} \langle 0, 0, 1 \rangle 8z \int_0^{a/2} \frac{dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}} \\
 &= \frac{2\lambda z}{\pi\epsilon_0} \hat{\mathbf{z}} \int_0^{a/2} \frac{dt}{\left(t^2 + \frac{a^2}{4} + z^2\right)^{3/2}}
 \end{aligned}$$

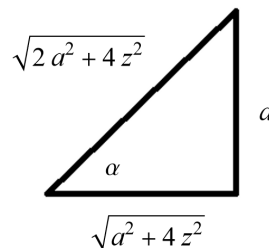
Make the following substitution.

$$\begin{aligned}
 t = \sqrt{\frac{a^2}{4} + z^2} \tan \theta &\quad \rightarrow \quad t^2 + \frac{a^2}{4} + z^2 = \left(\frac{a^2}{4} + z^2\right) (\tan^2 \theta + 1) = \left(\frac{a^2}{4} + z^2\right) \sec^2 \theta \\
 dt &= \sqrt{\frac{a^2}{4} + z^2} \sec^2 \theta d\theta
 \end{aligned}$$

As a result,

$$\begin{aligned}
 \mathbf{E} &= \frac{2\lambda z}{\pi\epsilon_0} \hat{\mathbf{z}} \int_{\tan^{-1}\left(\frac{0}{\sqrt{\frac{a^2}{4}+z^2}}\right)}^{\tan^{-1}\left(\frac{a/2}{\sqrt{\frac{a^2}{4}+z^2}}\right)} \frac{\sqrt{\frac{a^2}{4}+z^2} \sec^2 \theta d\theta}{\left[\left(\frac{a^2}{4}+z^2\right) \sec^2 \theta\right]^{3/2}} \\
 &= \frac{2\lambda z}{\pi\epsilon_0} \hat{\mathbf{z}} \int_0^{\tan^{-1}\left(\frac{a}{\sqrt{a^2+4z^2}}\right)} \frac{\sqrt{\frac{a^2}{4}+z^2} \sec^2 \theta d\theta}{\left(\frac{a^2}{4}+z^2\right)^{3/2} \sec^3 \theta} \\
 &= \frac{2\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{\frac{a^2}{4}+z^2} \int_0^{\tan^{-1}\left(\frac{a}{\sqrt{a^2+4z^2}}\right)} \cos \theta d\theta \\
 &= \frac{2\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{\frac{a^2}{4}+z^2} \sin \theta \Big|_0^{\tan^{-1}\left(\frac{a}{\sqrt{a^2+4z^2}}\right)} \\
 &= \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{a^2+4z^2} \sin \tan^{-1}\left(\frac{a}{\sqrt{a^2+4z^2}}\right).
 \end{aligned}$$

Draw the triangle implied by  $\alpha = \tan^{-1}\left(a/\sqrt{a^2+4z^2}\right)$  and use it to determine  $\sin \alpha$ .



$$\sin \alpha = \frac{a}{\sqrt{2a^2 + 4z^2}}$$

Therefore, the electric field at  $\mathbf{r} = \langle 0, 0, z \rangle$  is

$$\mathbf{E} = \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{a^2 + 4z^2} \left( \frac{a}{\sqrt{2a^2 + 4z^2}} \right).$$

Observe that

$$\begin{aligned}
 \lim_{a \rightarrow 0} \mathbf{E} &= \lim_{a \rightarrow 0} \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{a^2 + 4z^2} \left( \frac{a}{\sqrt{2a^2 + 4z^2}} \right) = \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{(0)^2 + 4z^2} \left( \frac{0}{\sqrt{2(0)^2 + 4z^2}} \right) = \mathbf{0} \\
 \lim_{z \rightarrow 0} \mathbf{E} &= \lim_{z \rightarrow 0} \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{a^2 + 4z^2} \left( \frac{a}{\sqrt{2a^2 + 4z^2}} \right) = \frac{8\lambda(0)}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{a^2 + 4(0)^2} \left( \frac{a}{\sqrt{2a^2 + 4(0)^2}} \right) = \mathbf{0}.
 \end{aligned}$$

In order to see what happens if  $z \gg a$ , rewrite the formula so that each term is a ratio of  $a$  and  $z$ ,  $z$  being in the denominator, and use the binomial theorem.

$$\begin{aligned}
 \mathbf{E} &= \frac{8\lambda z}{\pi\epsilon_0} \frac{\hat{\mathbf{z}}}{4z^2 \left(\frac{a^2}{4z^2} + 1\right)} \left( \frac{a}{2z\sqrt{\frac{a^2}{2z^2} + 1}} \right) \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \frac{1}{1 - \left(-\frac{a^2}{4z^2}\right)} \left(1 + \frac{a^2}{2z^2}\right)^{-1/2} \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left[ \sum_{k=0}^{\infty} \left(-\frac{a^2}{4z^2}\right)^k \right] \left[ \sum_{k=0}^{\infty} \frac{\Gamma\left(-\frac{1}{2} + 1\right)}{\Gamma(k+1)\Gamma\left(-\frac{1}{2} - k + 1\right)} \left(\frac{a^2}{2z^2}\right)^k \right] \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left[ \sum_{k=0}^{\infty} \left(-\frac{a^2}{4z^2}\right)^k \right] \left[ \sum_{k=0}^{\infty} \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(k+1)\Gamma\left(\frac{1}{2} - k\right)} \left(\frac{a^2}{2z^2}\right)^k \right] \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left[ \left(-\frac{a^2}{4z^2}\right)^0 + \left(-\frac{a^2}{4z^2}\right)^1 + \left(-\frac{a^2}{4z^2}\right)^2 + \dots \right] \\
 &\quad \times \left[ \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(1)\Gamma\left(\frac{1}{2}\right)} \left(\frac{a^2}{2z^2}\right)^0 + \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(2)\Gamma\left(-\frac{1}{2}\right)} \left(\frac{a^2}{2z^2}\right)^1 + \frac{\Gamma\left(\frac{1}{2}\right)}{\Gamma(3)\Gamma\left(-\frac{3}{2}\right)} \left(\frac{a^2}{2z^2}\right)^2 + \dots \right] \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left(1 - \frac{a^2}{4z^2} + \frac{a^4}{16z^4} - \dots\right) \left[1 - \frac{1}{2} \left(\frac{a^2}{2z^2}\right) + \frac{3}{8} \left(\frac{a^4}{4z^4}\right) - \dots\right] \\
 &= \frac{\lambda a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left(1 - \frac{a^2}{2z^2} + \frac{7a^4}{32z^4} - \dots\right) \\
 &= \left(\frac{Q}{4a}\right) \frac{a \hat{\mathbf{z}}}{\pi\epsilon_0 z^2} \left(1 - \frac{a^2}{2z^2} + \frac{7a^4}{32z^4} - \dots\right) \\
 &= \frac{Q \hat{\mathbf{z}}}{4\pi\epsilon_0 z^2} \left(1 - \frac{a^2}{2z^2} + \frac{7a^4}{32z^4} - \dots\right)
 \end{aligned}$$

If  $z \gg a$ , then  $a^2/z^2$  and all higher-order terms are so much smaller than 1 that they can be neglected.

$$\mathbf{E} \approx \frac{Q \hat{\mathbf{z}}}{4\pi\epsilon_0 z^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{z^2} \hat{\mathbf{z}}$$

The lesson is that far away from the square loop the electric field is the same as if it were a point charge.