

Problem 3.11

Find the momentum-space wave function, $\Phi(p, t)$, for a particle in the ground state of the harmonic oscillator. What is the probability (to two significant digits) that a measurement of p on a particle in this state would yield a value outside the classical range (for the same energy)? *Hint:* Look in a math table under “Normal Distribution” or “Error Function” for the numerical part—or use Mathematica.

Solution

The general formulas for the Fourier transform of a function $f(x)$ and its corresponding inverse Fourier transform are as follows.

$$\begin{cases} F(k) = \sqrt{\frac{|b|}{(2\pi)^{1-a}}} \int_{-\infty}^{\infty} e^{ibkx} f(x) dx \\ f(x) = \sqrt{\frac{|b|}{(2\pi)^{1+a}}} \int_{-\infty}^{\infty} e^{-ibkx} F(k) dk \end{cases}$$

The Fourier transform can be used to solve linear partial differential equations over the whole line. Any choice for a and b is acceptable, and how one chooses to define the Fourier transform really comes down to personal preference. In Chapter 2, for example, the Schrödinger equation was solved using $a = 0$ and $b = -1$.

$$\begin{cases} \mathcal{F}\{\Psi(x, t)\} = \tilde{\Psi}(k, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ikx} \Psi(x, t) dx \\ \mathcal{F}^{-1}\{\tilde{\Psi}(k, t)\} = \Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{ikx} \tilde{\Psi}(k, t) dk \end{cases}$$

One choice for a and b is special in quantum mechanics, though: $a = 0$ and $b = -1/\hbar$.

$$\begin{cases} \mathcal{F}\{\Psi(x, t)\} = \Phi(p, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \Psi(x, t) dx \\ \mathcal{F}^{-1}\{\Phi(p, t)\} = \Psi(x, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{ipx/\hbar} \Phi(p, t) dp \end{cases}$$

$\Psi(x, t)$ is the position-space wave function because $|\Psi(x, t)|^2$ represents the probability distribution for the particle's position. On the other hand, $\Phi(p, t)$ is the momentum-space wave function because $|\Phi(p, t)|^2$ represents the probability distribution for the particle's momentum. These formulas are a result of solving the eigenvalue problem for the momentum operator.

$$\hat{p}f(x) = pf(x)$$

$$-i\hbar \frac{d}{dx} f(x) = pf(x)$$

$$\frac{df}{dx} = \frac{ip}{\hbar} f(x)$$

$$f(x) = Ae^{ipx/\hbar}$$

This is a non-normalizable function, so the spectrum is continuous, meaning the continuous Dirac-analogs of Equations 3.10 and 3.11 on page 93 apply. Since \hat{p} is a hermitian operator, the eigenfunctions associated with the real, distinct eigenvalues are orthogonal.

$$\begin{aligned}\langle f' | f \rangle &= \int_{-\infty}^{\infty} (Ae^{ip'x/\hbar})^* (Ae^{ipx/\hbar}) dx = \int_{-\infty}^{\infty} (A^* e^{-ip'x/\hbar}) (Ae^{ipx/\hbar}) dx \\ &= |A|^2 \int_{-\infty}^{\infty} e^{i(p-p')x/\hbar} dx \\ &= A^2 \left[2\pi \delta \left(\frac{p-p'}{\hbar} \right) \right] \\ &= A^2 [2\pi |-\hbar| \delta(p' - p)] \\ &= 2\pi \hbar A^2 \delta(p' - p)\end{aligned}$$

Determine A by requiring the magnitude of the delta function to be 1.

$$2\pi \hbar A^2 = 1 \quad \rightarrow \quad A = \frac{1}{\sqrt{2\pi \hbar}}$$

Consequently,

$$f(x) = \frac{1}{\sqrt{2\pi \hbar}} e^{ipx/\hbar}.$$

\hat{p} is a hermitian operator, so any function in position-space, including the one we're most interested in, $\Psi(x, t)$, can be expressed as a linear combination of its eigenfunctions.

$$\Psi(x, t) = \int_{-\infty}^{\infty} B(p, t) \left(\frac{1}{\sqrt{2\pi \hbar}} e^{ipx/\hbar} \right) dp$$

By comparing this to the general formulas, we see that this is a very special inverse Fourier transform, one where $a = 0$ and $b = -1/\hbar$. The position-space wave function for a particle in the ground state of the harmonic oscillator potential is (see Problem 2.10)

$$\Psi(x, t) = \psi_0(x) e^{-iE_0 t/\hbar} = \left(\frac{m\omega}{\pi \hbar} \right)^{1/4} \exp \left(-\frac{m\omega}{2\hbar} x^2 \right) e^{-i\omega t/2}.$$

Take the Fourier transform of $\Psi(x, t)$ in order to get the momentum-space wave function.

$$\begin{aligned}\Phi(p, t) &= \mathcal{F}\{\Psi(x, t)\} \\ &= \frac{1}{\sqrt{2\pi \hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \Psi(x, t) dx \\ &= \frac{1}{\sqrt{2\pi \hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \left(\frac{m\omega}{\pi \hbar} \right)^{1/4} \exp \left(-\frac{m\omega}{2\hbar} x^2 \right) e^{-i\omega t/2} dx \\ &= \frac{e^{-i\omega t/2}}{\sqrt{2\pi \hbar}} \left(\frac{m\omega}{\pi \hbar} \right)^{1/4} \int_{-\infty}^{\infty} \exp \left(-\frac{m\omega}{2\hbar} x^2 - \frac{ip}{\hbar} x \right) dx\end{aligned}$$

Complete the square in the exponent.

$$\begin{aligned}\Phi(p, t) &= \frac{e^{-i\omega t/2}}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \int_{-\infty}^{\infty} \exp\left[-\frac{m\omega}{2\hbar} \left(x^2 + \frac{2ip}{m\omega}x\right)\right] dx \\ &= \frac{e^{-i\omega t/2}}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \int_{-\infty}^{\infty} \exp\left[-\frac{m\omega}{2\hbar} \left(x^2 + \frac{2ip}{m\omega}x - \frac{p^2}{m^2\omega^2}\right)\right] \exp\left[-\frac{m\omega}{2\hbar} \left(\frac{p^2}{m^2\omega^2}\right)\right] dx \\ &= \frac{e^{-i\omega t/2}}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \exp\left(-\frac{p^2}{2\hbar m\omega}\right) \int_{-\infty}^{\infty} \exp\left[-\frac{m\omega}{2\hbar} \left(x + \frac{ip}{m\omega}\right)^2\right] dx\end{aligned}$$

Make the following substitution.

$$\begin{aligned}u &= \sqrt{\frac{m\omega}{2\hbar}} \left(x + \frac{ip}{m\omega}\right) \\ du &= \sqrt{\frac{m\omega}{2\hbar}} dx \quad \rightarrow \quad dx = \sqrt{\frac{2\hbar}{m\omega}} du\end{aligned}$$

As a result,

$$\begin{aligned}\Phi(p, t) &= \frac{e^{-i\omega t/2}}{\sqrt{2\pi\hbar}} \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \exp\left(-\frac{p^2}{2\hbar m\omega}\right) \int_{-\infty}^{\infty} e^{-u^2} \sqrt{\frac{2\hbar}{m\omega}} du \\ &= \frac{e^{-i\omega t/2}}{\sqrt{\pi}} \frac{1}{(\pi\hbar m\omega)^{1/4}} \exp\left(-\frac{p^2}{2\hbar m\omega}\right) \underbrace{\int_{-\infty}^{\infty} e^{-u^2} du}_{=\sqrt{\pi}}\end{aligned}$$

$$\boxed{\Phi(p, t) = \frac{e^{-i\omega t/2}}{(\pi\hbar m\omega)^{1/4}} \exp\left(-\frac{p^2}{2\hbar m\omega}\right).}$$

A classical particle with mass m and energy $E \geq 0$ has momentum

$$\begin{aligned}E &\geq \frac{p^2}{2m} \\ p^2 &\leq 2mE \\ |p| &\leq \sqrt{2mE} \\ -\sqrt{2mE} &\leq p \leq \sqrt{2mE}.\end{aligned}$$

Assuming it has the energy of the harmonic-oscillator ground state, $E = E_0 = \hbar\omega/2$,

$$-\sqrt{\hbar m\omega} \leq p \leq \sqrt{\hbar m\omega}.$$

The probability of measuring p outside of this range is

$$P = \int_{-\infty}^{-\sqrt{\hbar m\omega}} |\Phi(p, t)|^2 dp + \int_{\sqrt{\hbar m\omega}}^{\infty} |\Phi(p, t)|^2 dp$$

because $|\Phi(p, t)|^2$ is the probability distribution for the particle's momentum.

$$\begin{aligned}
 P &= \int_{-\infty}^{-\sqrt{\hbar m \omega}} \left| \frac{e^{-i\omega t/2}}{(\pi \hbar m \omega)^{1/4}} \exp\left(-\frac{p^2}{2\hbar m \omega}\right) \right|^2 dp + \int_{\sqrt{\hbar m \omega}}^{\infty} \left| \frac{e^{-i\omega t/2}}{(\pi \hbar m \omega)^{1/4}} \exp\left(-\frac{p^2}{2\hbar m \omega}\right) \right|^2 dp \\
 &= \int_{-\infty}^{-\sqrt{\hbar m \omega}} \frac{1}{\sqrt{\pi \hbar m \omega}} \exp\left(-\frac{p^2}{\hbar m \omega}\right) dp + \int_{\sqrt{\hbar m \omega}}^{\infty} \frac{1}{\sqrt{\pi \hbar m \omega}} \exp\left(-\frac{p^2}{\hbar m \omega}\right) dp \\
 &= \frac{1}{\sqrt{\pi \hbar m \omega}} \left[\int_{-\infty}^{-\sqrt{\hbar m \omega}} \exp\left(-\frac{p^2}{\hbar m \omega}\right) dp + \int_{\sqrt{\hbar m \omega}}^{\infty} \exp\left(-\frac{p^2}{\hbar m \omega}\right) dp \right]
 \end{aligned}$$

Make the following substitutions.

$$\begin{aligned}
 v &= -\frac{p}{\sqrt{\hbar m \omega}} & w &= \frac{p}{\sqrt{\hbar m \omega}} \\
 dv &= -\frac{dp}{\sqrt{\hbar m \omega}} \quad \rightarrow \quad dp = -\sqrt{\hbar m \omega} dv & dw &= \frac{dp}{\sqrt{\hbar m \omega}} \quad \rightarrow \quad dp = \sqrt{\hbar m \omega} dw
 \end{aligned}$$

As a result,

$$\begin{aligned}
 P &= \frac{1}{\sqrt{\pi \hbar m \omega}} \left[\int_{\infty}^1 e^{-v^2} (-\sqrt{\hbar m \omega} dv) + \int_1^{\infty} e^{-w^2} (\sqrt{\hbar m \omega} dw) \right] \\
 &= \frac{1}{\sqrt{\pi}} \left(\int_1^{\infty} e^{-v^2} dv + \int_1^{\infty} e^{-w^2} dw \right) \\
 &= \frac{2}{\sqrt{\pi}} \int_1^{\infty} e^{-v^2} dv \\
 &\approx \frac{2}{\sqrt{\pi}} (0.139403) \\
 &\approx 0.16.
 \end{aligned}$$