

Problem 3.34

- (a) Find the momentum-space wave function $\Phi_n(p, t)$ for the n th stationary state of the infinite square well.
- (b) Find the probability density $|\Phi_n(p, t)|^2$. Graph this function, for $n = 1$, $n = 2$, $n = 5$, and $n = 10$. What are the most probable values of p , for large n ? Is this what you would have expected?⁴⁰ Compare your answer to Problem 3.10.
- (c) Use $\Phi_n(p, t)$ to calculate the expectation value of p^2 , in the n th state. Compare your answer to Problem 2.4.

Solution

The position-space wave function of the infinite square well is

$$\Psi_n(x, t) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right), \quad 0 \leq x \leq a,$$

where $n = 1, 2, \dots$, and zero elsewhere. To find the momentum-space wave function, take the Fourier transform of $\Psi_n(x, t)$.

$$\begin{aligned} \Phi_n(p, t) &= \mathcal{F}\{\Psi_n(x, t)\} \\ &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \Psi_n(x, t) dx \\ &= \frac{1}{\sqrt{2\pi\hbar}} \int_0^a e^{-ipx/\hbar} \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) dx \\ &= \frac{1}{\sqrt{\pi\hbar a}} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) \int_0^a e^{-ipx/\hbar} \sin \frac{n\pi x}{a} dx \\ &= \frac{1}{\sqrt{\pi\hbar a}} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) \int_0^a e^{-ipx/\hbar} \left(\frac{e^{in\pi x/a} - e^{-in\pi x/a}}{2i}\right) dx \\ &= \frac{1}{2i\sqrt{\pi\hbar a}} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) \int_0^a \left\{ \exp\left[i\left(-\frac{p}{\hbar} + \frac{n\pi}{a}\right)x\right] - \exp\left[-i\left(\frac{p}{\hbar} + \frac{n\pi}{a}\right)x\right] \right\} dx \\ &= \frac{1}{2i\sqrt{\pi\hbar a}} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) \int_0^a \left[\exp\left(i \frac{n\pi\hbar - pa}{\hbar a} x\right) - \exp\left(-i \frac{n\pi\hbar + pa}{\hbar a} x\right) \right] dx \\ &= \frac{1}{2i\sqrt{\pi\hbar a}} \exp\left(-i \frac{n^2 \pi^2 \hbar}{2ma^2} t\right) \left[\int_0^a \exp\left(i \frac{n\pi\hbar - pa}{\hbar a} x\right) dx - \int_0^a \exp\left(-i \frac{n\pi\hbar + pa}{\hbar a} x\right) dx \right] \end{aligned}$$

⁴⁰See F. L. Markley, *Am. J. Phys.* **40**, 1545 (1972).

Evaluate the integrals and simplify the formula.

$$\begin{aligned}
\Phi_n(p, t) &= \frac{1}{2i\sqrt{\pi\hbar a}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left[\frac{\hbar a}{i(n\pi\hbar - pa)} \exp\left(i\frac{n\pi\hbar - pa}{\hbar}x\right) \Big|_0^a - \frac{\hbar a}{-i(n\pi\hbar + pa)} \exp\left(-i\frac{n\pi\hbar + pa}{\hbar}x\right) \Big|_0^a \right] \\
&= \frac{1}{2i^2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left\{ \frac{1}{n\pi\hbar - pa} \left[\exp\left(i\frac{n\pi\hbar - pa}{\hbar}\right) - 1 \right] + \frac{1}{n\pi\hbar + pa} \left[\exp\left(-i\frac{n\pi\hbar + pa}{\hbar}\right) - 1 \right] \right\} \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left[\frac{\exp\left(i\frac{n\pi\hbar - pa}{\hbar}\right)}{n\pi\hbar - pa} - \frac{1}{n\pi\hbar - pa} + \frac{\exp\left(-i\frac{n\pi\hbar + pa}{\hbar}\right)}{n\pi\hbar + pa} - \frac{1}{n\pi\hbar + pa} \right] \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left[\frac{\exp\left(i\frac{n\pi\hbar - pa}{\hbar}\right)(n\pi\hbar + pa) - (n\pi\hbar + pa) + \exp\left(-i\frac{n\pi\hbar + pa}{\hbar}\right)(n\pi\hbar - pa) - (n\pi\hbar - pa)}{(n\pi\hbar - pa)(n\pi\hbar + pa)} \right] \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left\{ \frac{n\pi\hbar \left[\exp\left(i\frac{n\pi\hbar - pa}{\hbar}\right) + \exp\left(-i\frac{n\pi\hbar + pa}{\hbar}\right) \right] + pa \left[\exp\left(i\frac{n\pi\hbar - pa}{\hbar}\right) - \exp\left(-i\frac{n\pi\hbar + pa}{\hbar}\right) \right] - 2n\pi\hbar}{n^2\pi^2\hbar^2 - p^2a^2} \right\} \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left[\frac{n\pi\hbar e^{-ipa/\hbar} (e^{in\pi} + e^{-in\pi}) + pa e^{-ipa/\hbar} (e^{in\pi} - e^{-in\pi}) - 2n\pi\hbar}{n^2\pi^2\hbar^2 - p^2a^2} \right] \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left[\frac{n\pi\hbar e^{-ipa/\hbar} (2 \cos n\pi) + pa e^{-ipa/\hbar} (2i \sin n\pi) - 2n\pi\hbar}{n^2\pi^2\hbar^2 - p^2a^2} \right] \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left\{ \frac{n\pi\hbar e^{-ipa/\hbar} [2(-1)^n] + pa e^{-ipa/\hbar} [2i(0)] - 2n\pi\hbar}{n^2\pi^2\hbar^2 - p^2a^2} \right\} \\
&= -\frac{1}{2} \sqrt{\frac{\hbar a}{\pi}} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \left\{ \frac{2n\pi\hbar [e^{-ipa/\hbar}(-1)^n - 1]}{n^2\pi^2\hbar^2 - p^2a^2} \right\} \\
&= -\sqrt{\pi\hbar^3 a} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \frac{ne^{-ipa/(2\hbar)} [e^{ipa/(2\hbar)} - (-1)^n e^{-ipa/(2\hbar)}]}{p^2a^2 - n^2\pi^2\hbar^2}
\end{aligned}$$

Simplify the formula even further by considering the cases where n is odd and even separately.

$$\begin{aligned}\Phi_n(p, t) &= \begin{cases} -\sqrt{\pi\hbar^3 a} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \frac{ne^{-ipa/(2\hbar)}[e^{ipa/(2\hbar)} + e^{-ipa/(2\hbar)}]}{p^2a^2 - n^2\pi^2\hbar^2} & \text{if } n = 2k - 1 \\ -\sqrt{\pi\hbar^3 a} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \frac{ne^{-ipa/(2\hbar)}[e^{ipa/(2\hbar)} - e^{-ipa/(2\hbar)}]}{p^2a^2 - n^2\pi^2\hbar^2} & \text{if } n = 2k \end{cases} \\ &= \begin{cases} -\sqrt{\pi\hbar^3 a} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \frac{ne^{-ipa/(2\hbar)}(2\cos\frac{pa}{2\hbar})}{p^2a^2 - n^2\pi^2\hbar^2} & \text{if } n = 2k - 1 \\ -\sqrt{\pi\hbar^3 a} \exp\left(-i\frac{n^2\pi^2\hbar}{2ma^2}t\right) \frac{ne^{-ipa/(2\hbar)}(2i\sin\frac{pa}{2\hbar})}{p^2a^2 - n^2\pi^2\hbar^2} & \text{if } n = 2k \end{cases}\end{aligned}$$

Now calculate $|\Phi_n(p, t)|^2$, the probability distribution for the particle's momentum in the n th stationary state of the infinite square well.

$$\begin{aligned}|\Phi_n(p, t)|^2 &= \Phi_n^*(p, t)\Phi_n(p, t) \\ &= \begin{cases} \pi\hbar^3 a \frac{4n^2 \cos^2 \frac{pa}{2\hbar}}{(p^2a^2 - n^2\pi^2\hbar^2)^2} & \text{if } n = 2k - 1 \\ \pi\hbar^3 a \frac{4n^2 \sin^2 \frac{pa}{2\hbar}}{(p^2a^2 - n^2\pi^2\hbar^2)^2} & \text{if } n = 2k \end{cases} \\ &= \begin{cases} \frac{\pi a}{\hbar} \frac{4n^2 \cos^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar}\right)^2 - n^2\pi^2\right]^2} & \text{if } n = 2k - 1 \\ \frac{\pi a}{\hbar} \frac{4n^2 \sin^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar}\right)^2 - n^2\pi^2\right]^2} & \text{if } n = 2k \end{cases}\end{aligned}$$

Check that it's normalized.

$$\int_{-\infty}^{\infty} |\Phi_n(p, t)|^2 dp = \begin{cases} \frac{4\pi an^2}{\hbar} \int_{-\infty}^{\infty} \frac{\cos^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar}\right)^2 - n^2\pi^2\right]^2} dp & \text{if } n = 2k - 1 \\ \frac{4\pi an^2}{\hbar} \int_{-\infty}^{\infty} \frac{\sin^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar}\right)^2 - n^2\pi^2\right]^2} dp & \text{if } n = 2k \end{cases}$$

Make the following substitution.

$$\begin{aligned}\frac{pa}{\hbar} = n\pi u &\quad \rightarrow \quad \frac{pa}{2\hbar} = \frac{n\pi u}{2} \\ dp = \frac{n\pi\hbar}{a} du\end{aligned}$$

Consequently, using partial fraction decomposition,

$$\begin{aligned}
 \int_{-\infty}^{\infty} |\Phi_n(p, t)|^2 dp &= \begin{cases} \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(n^2\pi^2 u^2 - n^2\pi^2)^2} \left(\frac{n\pi\hbar}{a} du \right) & \text{if } n = 2k - 1 \\ \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(n^2\pi^2 u^2 - n^2\pi^2)^2} \left(\frac{n\pi\hbar}{a} du \right) & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{4}{\pi^2 n} \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(u^2 - 1)^2} du & \text{if } n = 2k - 1 \\ \frac{4}{\pi^2 n} \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(u^2 - 1)^2} du & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{4}{\pi^2 n} \int_{-\infty}^{\infty} \left[\frac{1}{4(u+1)} + \frac{1}{4(u+1)^2} - \frac{1}{4(u-1)} + \frac{1}{4(u-1)^2} \right] \cos^2 \frac{n\pi u}{2} du & \text{if } n = 2k - 1 \\ \frac{4}{\pi^2 n} \int_{-\infty}^{\infty} \left[\frac{1}{4(u+1)} + \frac{1}{4(u+1)^2} - \frac{1}{4(u-1)} + \frac{1}{4(u-1)^2} \right] \sin^2 \frac{n\pi u}{2} du & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{1}{\pi^2 n} \left[\int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{u+1} du + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(u+1)^2} du - \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{u-1} du + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(u-1)^2} du \right] & \text{if } n = 2k - 1 \\ \frac{1}{\pi^2 n} \left[\int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{u+1} du + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(u+1)^2} du - \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{u-1} du + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(u-1)^2} du \right] & \text{if } n = 2k \end{cases} .
 \end{aligned}$$

Make the substitutions, $v = u + 1$ and $w = u - 1$.

$$\int_{-\infty}^{\infty} |\Phi_n(p, t)|^2 dp = \begin{cases} \frac{1}{\pi^2 n} \left[\int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(v-1)}{2}}{v} dv + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(v-1)}{2}}{v^2} dv - \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(w+1)}{2}}{w} dw + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(w+1)}{2}}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{1}{\pi^2 n} \left[\int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(v-1)}{2}}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(v-1)}{2}}{v^2} dv - \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(w+1)}{2}}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(w+1)}{2}}{w^2} dw \right] & \text{if } n = 2k \end{cases}$$

$$= \begin{cases} \frac{1}{\pi^2(2k-1)} \left[\int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi v - \frac{\pi v}{2} - k\pi + \frac{\pi}{2} \right)}{v} dv + \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi v - \frac{\pi v}{2} - k\pi + \frac{\pi}{2} \right)}{v^2} dv \right. \\ \left. - \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi w - \frac{\pi w}{2} + k\pi - \frac{\pi}{2} \right)}{w} dw + \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi w - \frac{\pi w}{2} + k\pi - \frac{\pi}{2} \right)}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{1}{\pi^2(2k)} \left[\int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v^2} dv - \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w^2} dw \right] & \text{if } n = 2k \end{cases}$$

$$= \begin{cases} \frac{1}{\pi^2(2k-1)} \left[\int_{-\infty}^{\infty} \frac{[-\sin \left(k\pi v - \frac{\pi v}{2} - k\pi \right)]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[-\sin \left(k\pi v - \frac{\pi v}{2} - k\pi \right)]^2}{v^2} dv \right. \\ \left. - \int_{-\infty}^{\infty} \frac{[\sin \left(k\pi w - \frac{\pi w}{2} + k\pi \right)]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[\sin \left(k\pi w - \frac{\pi w}{2} + k\pi \right)]^2}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{1}{\pi^2(2k)} \left[\int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v^2} dv - \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w^2} dw \right] & \text{if } n = 2k \end{cases}$$

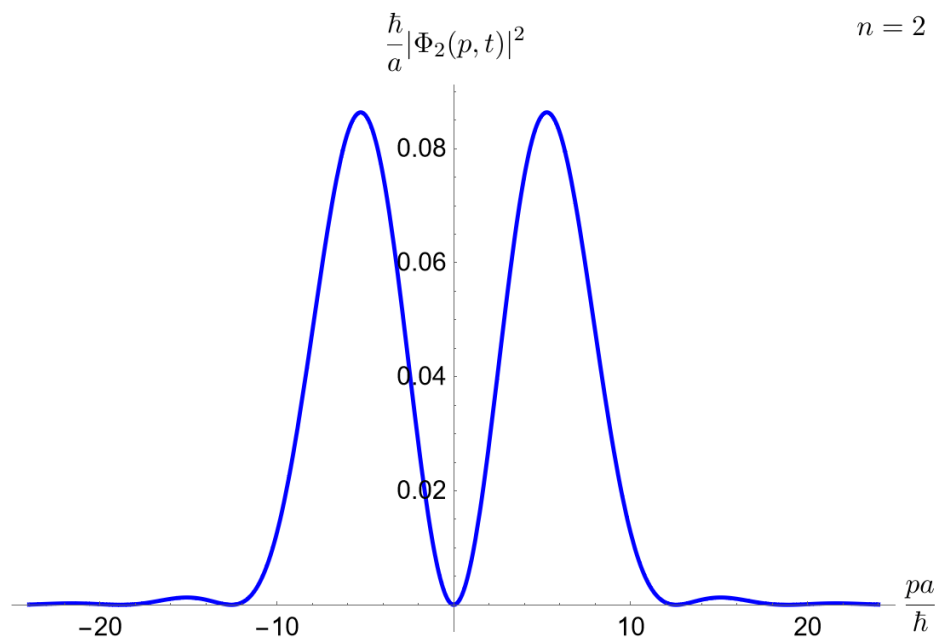
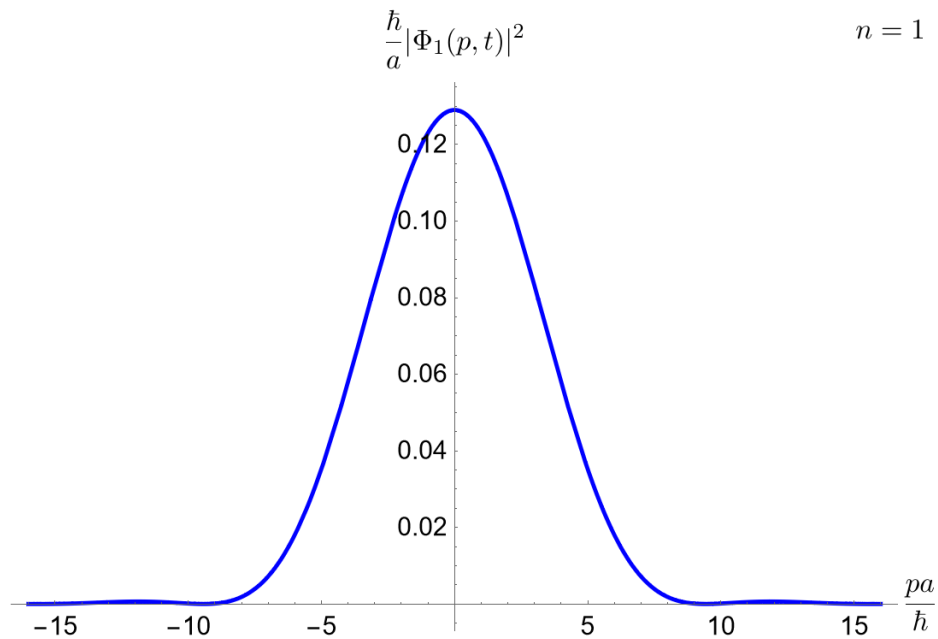
Continue the simplification.

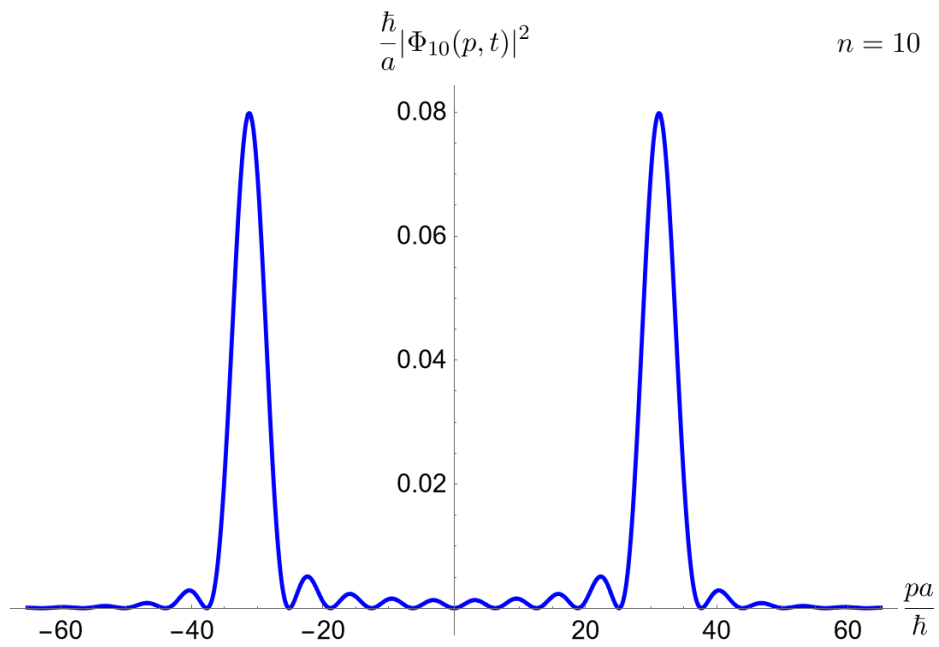
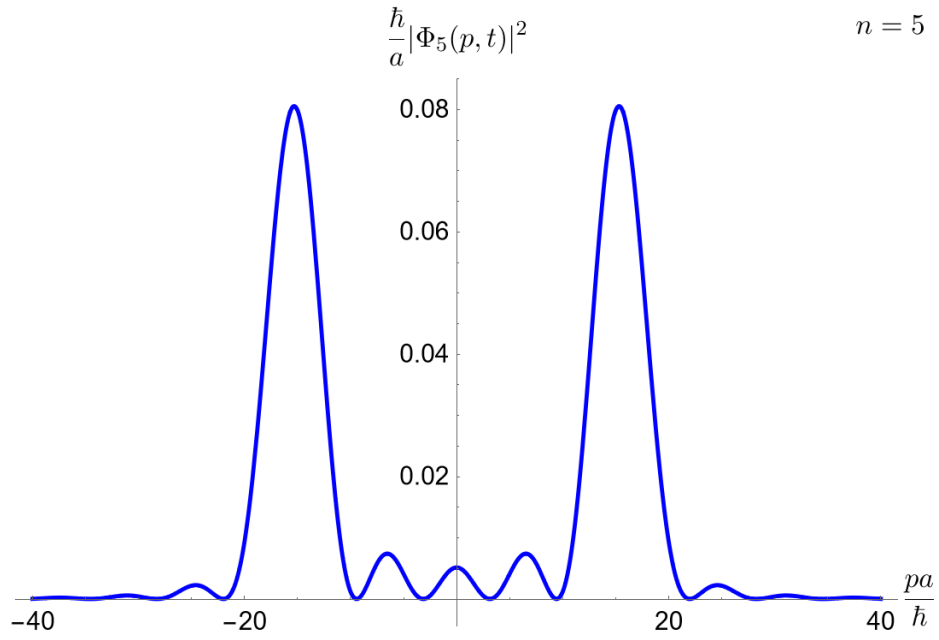
$$\begin{aligned}
 \int_{-\infty}^{\infty} |\Phi_n(p, t)|^2 dp &= \begin{cases} \frac{1}{\pi^2(2k-1)} \left[\int_{-\infty}^{\infty} \frac{[-(-1)^k \sin(k\pi v - \frac{\pi v}{2})]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[-(-1)^k \sin(k\pi v - \frac{\pi v}{2})]^2}{v^2} dv \right. \\ \qquad \qquad \qquad \left. - \int_{-\infty}^{\infty} \frac{[(-1)^k \sin(k\pi w - \frac{\pi w}{2})]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin(k\pi w - \frac{\pi w}{2})]^2}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \\ \frac{1}{\pi^2(2k)} \left[\int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi v]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi v]^2}{v^2} dv \right. \\ \qquad \qquad \qquad \left. - \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi w]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi w]^2}{w^2} dw \right] & \text{if } n = 2k \end{cases} \\
 \\
 &= \begin{cases} \frac{1}{\pi^2(2k-1)} \left[\int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - \frac{\pi v}{2})}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - \frac{\pi v}{2})}{v^2} dv \right. \\ \qquad \qquad \qquad \left. - \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w - \frac{\pi w}{2})}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w - \frac{\pi w}{2})}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \\ \frac{1}{\pi^2(2k)} \left(\int_{-\infty}^{\infty} \frac{\sin^2 k\pi v}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2 k\pi v}{v^2} dv - \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w^2} dw \right) & \text{if } n = 2k \end{cases} \\
 \\
 &= \begin{cases} \frac{2}{\pi^2(2k-1)} \int_{-\infty}^{\infty} \frac{\sin^2(k - \frac{1}{2})\pi w}{w^2} dw & \text{if } n = 2k - 1 \\ \\ \frac{2}{\pi^2(2k)} \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w^2} dw & \text{if } n = 2k \end{cases}
 \end{aligned}$$

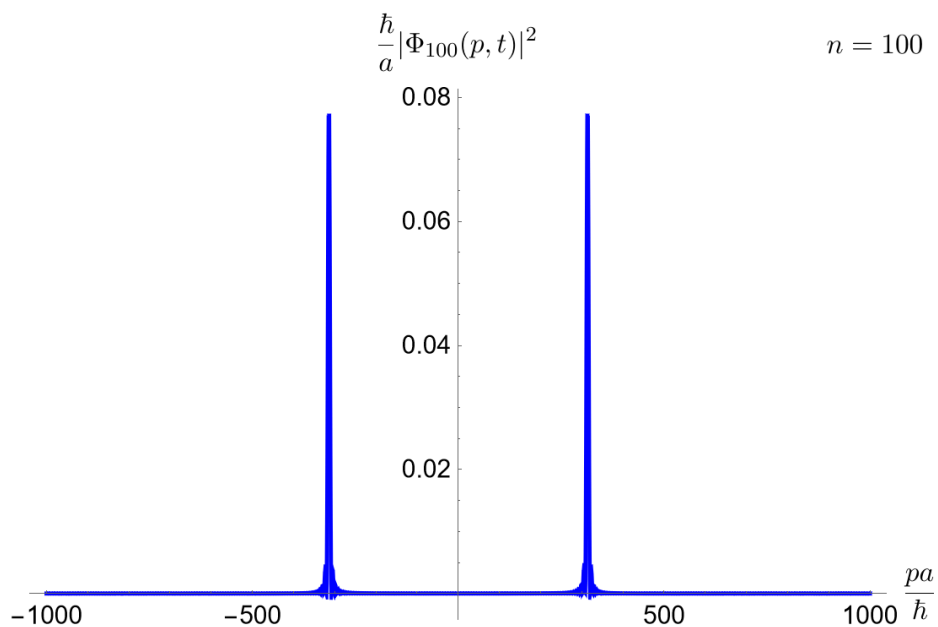
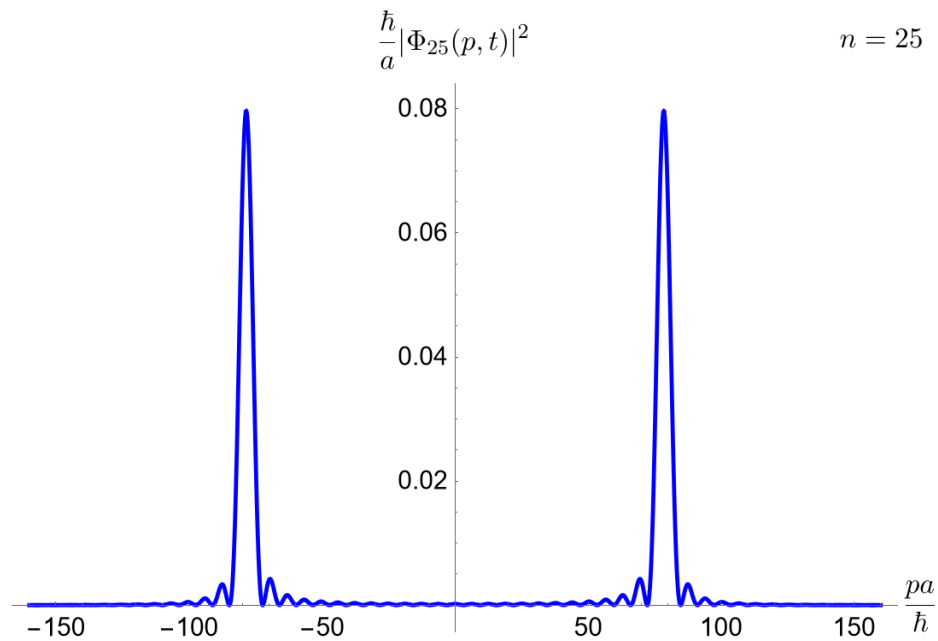
Evaluate the integrals and finally confirm that the momentum-space wave function is normalized.

$$\int_{-\infty}^{\infty} |\Phi_n(p, t)|^2 dp = \begin{cases} \frac{2}{\pi^2(2k-1)} \left[\left(k - \frac{1}{2}\right) \pi^2 \right] & \text{if } n = 2k - 1 \\ \frac{2}{\pi^2(2k)} (k\pi^2) & \text{if } n = 2k \end{cases} = 1$$

Below are plots of $(\hbar/a)|\Phi_n(p, t)|^2$ versus pa/\hbar for various values of n .







The most probable values of p occur where the absolute maxima are in these graphs. Differentiate $|\Phi(p, t)|^2$ with respect to p and set it equal to zero.

$$\frac{\partial}{\partial p} |\Phi(p, t)|^2 = \begin{cases} \frac{2a^2 n^2 \pi \hbar^2 [-(p^2 a^2 - n^2 \pi^2 \hbar^2) \sin \frac{pa}{\hbar} - 4pa\hbar (1 + \cos \frac{pa}{\hbar})]}{(p^2 a^2 - n^2 \pi^2 \hbar^2)^3} & \text{if } n = 2k - 1 \\ \frac{2a^2 n^2 \pi \hbar^2 [(p^2 a^2 - n^2 \pi^2 \hbar^2) \sin \frac{pa}{\hbar} - 4pa\hbar (1 - \cos \frac{pa}{\hbar})]}{(p^2 a^2 - n^2 \pi^2 \hbar^2)^3} & \text{if } n = 2k \end{cases} = 0$$

The derivative vanishes when the quantities in square brackets do.

$$\begin{cases} -(p^2 a^2 - n^2 \pi^2 \hbar^2) \sin \frac{pa}{\hbar} - 4pa\hbar \left(1 + \cos \frac{pa}{\hbar}\right) = 0 & \text{if } n = 2k - 1 \\ (p^2 a^2 - n^2 \pi^2 \hbar^2) \sin \frac{pa}{\hbar} - 4pa\hbar \left(1 - \cos \frac{pa}{\hbar}\right) = 0 & \text{if } n = 2k \end{cases}$$

Divide both sides of each equation by \hbar^2 .

$$\begin{cases} -\left(\frac{p^2 a^2}{\hbar^2} - n^2 \pi^2\right) \sin \frac{pa}{\hbar} - \frac{4pa}{\hbar} \left(1 + \cos \frac{pa}{\hbar}\right) = 0 & \text{if } n = 2k - 1 \\ \left(\frac{p^2 a^2}{\hbar^2} - n^2 \pi^2\right) \sin \frac{pa}{\hbar} - \frac{4pa}{\hbar} \left(1 - \cos \frac{pa}{\hbar}\right) = 0 & \text{if } n = 2k \end{cases}$$

Set $z = pa/\hbar$.

$$\begin{cases} -(z^2 - n^2 \pi^2) \sin z - 4z(1 + \cos z) = 0 & \text{if } n = 2k - 1 \\ (z^2 - n^2 \pi^2) \sin z - 4z(1 - \cos z) = 0 & \text{if } n = 2k \end{cases}$$

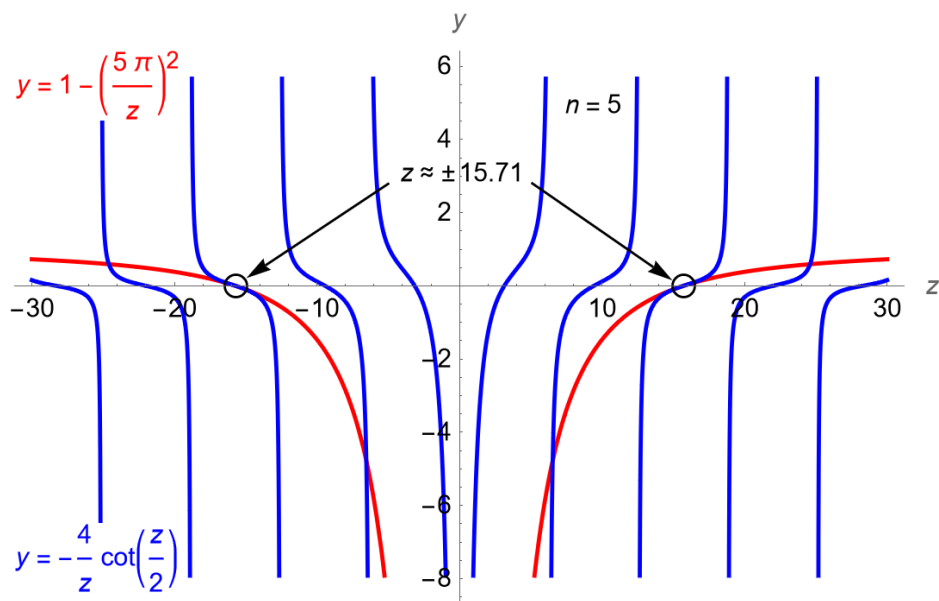
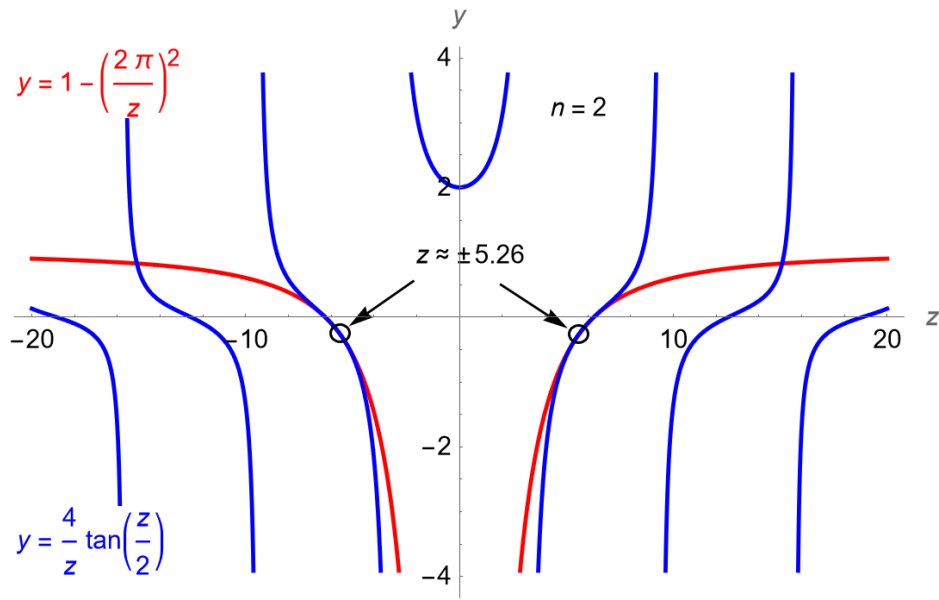
Rearrange the terms, assuming $z \neq n\pi$.

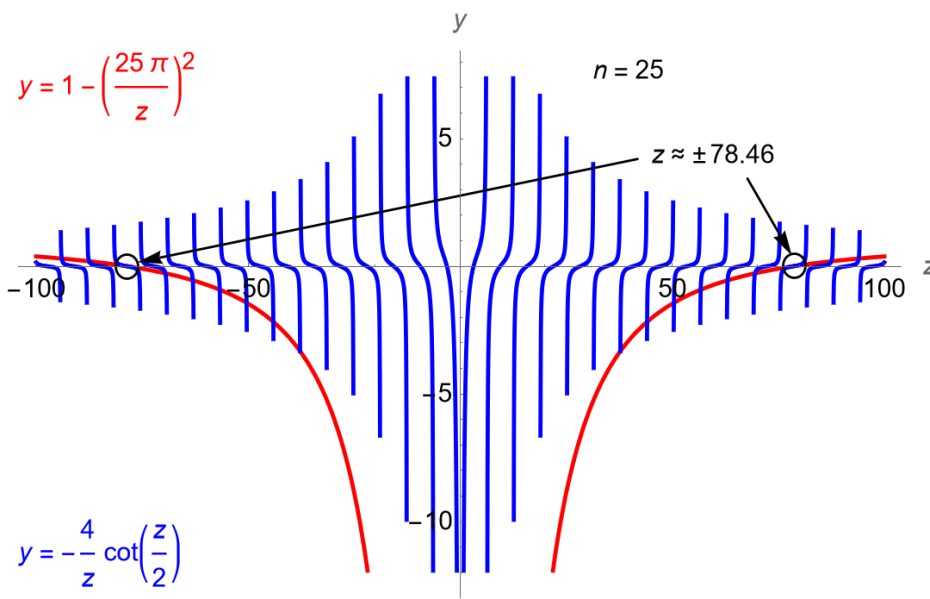
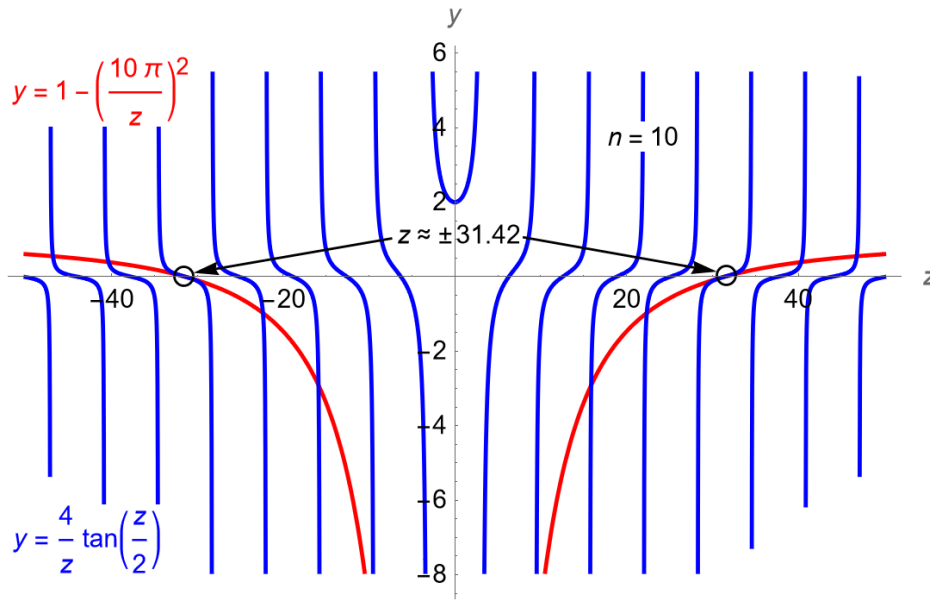
$$\begin{cases} z^2 - n^2 \pi^2 = -4z \left(\frac{1 + \cos z}{\sin z}\right) & \text{if } n = 2k - 1 \\ z^2 - n^2 \pi^2 = 4z \left(\frac{1 - \cos z}{\sin z}\right) & \text{if } n = 2k \end{cases}$$

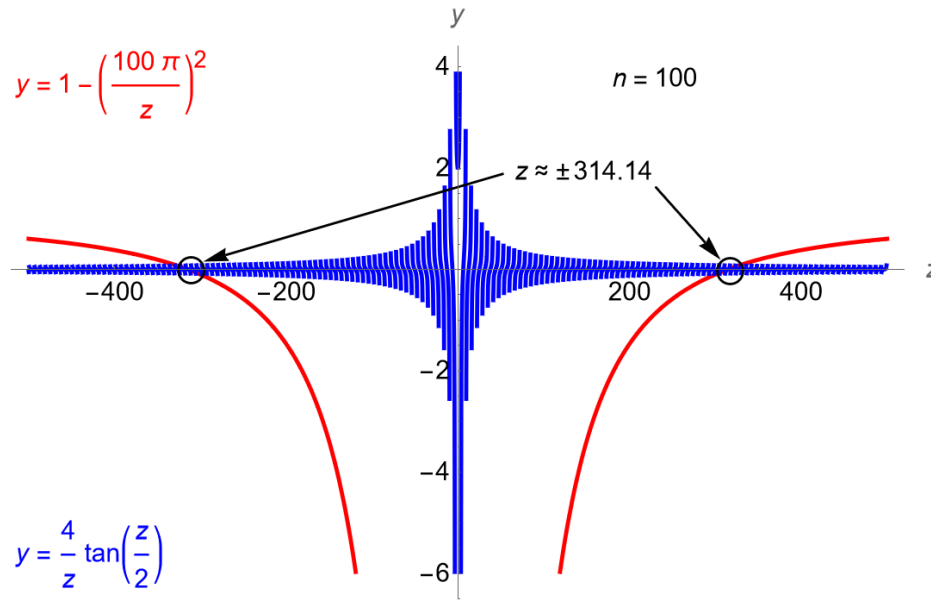
Divide both sides by z^2 and substitute the appropriate half-angle formulas.

$$\begin{cases} 1 - \left(\frac{n\pi}{z}\right)^2 = -\frac{4}{z} \cot \frac{z}{2} & \text{if } n = 2k - 1 \\ 1 - \left(\frac{n\pi}{z}\right)^2 = \frac{4}{z} \tan \frac{z}{2} & \text{if } n = 2k \end{cases}$$

Graph the functions on both sides versus z for the values of n used previously and label the intersections where the absolute maxima are. For $n = 1$, the absolute maximum is at $z = 0$.







Observe that the intersections where the absolute maxima occur roughly on the z -axis when $n \geq 5$.

$$\begin{cases} 1 - \left(\frac{n\pi}{z}\right)^2 = -\frac{4}{z} \cot \frac{z}{2} \approx 0 & \text{if } n = 2k - 1 \\ 1 - \left(\frac{n\pi}{z}\right)^2 = \frac{4}{z} \tan \frac{z}{2} \approx 0 & \text{if } n = 2k \end{cases}$$

Solving for z yields $z \approx \pm n\pi$. Therefore, the most probable values of p for large n are

$$\frac{pa}{\hbar} \approx \pm n\pi \quad \rightarrow \quad p \approx \pm \frac{n\pi\hbar}{a}.$$

The Hamiltonian for a particle in the infinite square well is equal to the total mechanical energy.

$$H = T + V = \frac{p^2}{2m} + 0 = \frac{p^2}{2m} = E$$

Solve for p .

$$p = \pm\sqrt{2mE}$$

Plug in the formula for the energy of the n th stationary state.

$$p = \pm\sqrt{2m \left(\frac{n^2\pi^2\hbar^2}{2ma^2}\right)} = \pm\sqrt{\frac{n^2\pi^2\hbar^2}{a^2}} = \pm\frac{n\pi\hbar}{a}$$

In the limit as $n \rightarrow \infty$, then, the particle has the momenta predicted by classical mechanics.

The aim now is to calculate the expectation value of p^2 of the n th stationary state at time t using the momentum-space wave function.

$$\begin{aligned}
 \langle p^2 \rangle &= \langle \Phi | \hat{p}^2 | \Phi \rangle \\
 &= \int_{-\infty}^{\infty} \Phi_n^*(p, t) p^2 \Phi_n(p, t) dp \\
 &= \int_{-\infty}^{\infty} p^2 |\Phi_n(p, t)|^2 dp \\
 &= \begin{cases} \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{p^2 \cos^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar} \right)^2 - n^2 \pi^2 \right]^2} dp & \text{if } n = 2k - 1 \\ \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{p^2 \sin^2 \frac{pa}{2\hbar}}{\left[\left(\frac{pa}{\hbar} \right)^2 - n^2 \pi^2 \right]^2} dp & \text{if } n = 2k \end{cases}
 \end{aligned}$$

Make the following substitution.

$$\begin{aligned}
 \frac{pa}{\hbar} = n\pi u &\quad \rightarrow \quad \frac{pa}{2\hbar} = \frac{n\pi u}{2} &\quad \rightarrow \quad p = \frac{n\pi\hbar}{a} u \\
 dp = \frac{n\pi\hbar}{a} du &
 \end{aligned}$$

Consequently, using partial fraction decomposition,

$$\begin{aligned}
 \langle p^2 \rangle &= \begin{cases} \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{\left(\frac{n\pi\hbar}{a} u \right)^2 \cos^2 \frac{n\pi u}{2}}{(n^2 \pi^2 u^2 - n^2 \pi^2)^2} \left(\frac{n\pi\hbar}{a} du \right) & \text{if } n = 2k - 1 \\ \frac{4\pi a n^2}{\hbar} \int_{-\infty}^{\infty} \frac{\left(\frac{n\pi\hbar}{a} u \right)^2 \sin^2 \frac{n\pi u}{2}}{(n^2 \pi^2 u^2 - n^2 \pi^2)^2} \left(\frac{n\pi\hbar}{a} du \right) & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{4n\hbar^2}{a^2} \int_{-\infty}^{\infty} \frac{u^2 \cos^2 \frac{n\pi u}{2}}{(u^2 - 1)^2} du & \text{if } n = 2k - 1 \\ \frac{4n\hbar^2}{a^2} \int_{-\infty}^{\infty} \frac{u^2 \sin^2 \frac{n\pi u}{2}}{(u^2 - 1)^2} du & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{4n\hbar^2}{a^2} \int_{-\infty}^{\infty} \left[-\frac{1}{4(u+1)} + \frac{1}{4(u+1)^2} + \frac{1}{4(u-1)} + \frac{1}{4(u-1)^2} \right] \cos^2 \frac{n\pi u}{2} du & \text{if } n = 2k - 1 \\ \frac{4n\hbar^2}{a^2} \int_{-\infty}^{\infty} \left[-\frac{1}{4(u+1)} + \frac{1}{4(u+1)^2} + \frac{1}{4(u-1)} + \frac{1}{4(u-1)^2} \right] \sin^2 \frac{n\pi u}{2} du & \text{if } n = 2k \end{cases}
 \end{aligned}$$

Split up the integrals and then make the substitutions, $v = u + 1$ and $w = u - 1$.

$$\begin{aligned}
 \langle p^2 \rangle &= \begin{cases} \frac{n\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{u+1} du + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(u+1)^2} du + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{u-1} du + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi u}{2}}{(u-1)^2} du \right] & \text{if } n = 2k - 1 \\ \frac{n\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{u+1} du + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(u+1)^2} du + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{u-1} du + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi u}{2}}{(u-1)^2} du \right] & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{n\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(v-1)}{2}}{v} dv + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(v-1)}{2}}{v^2} dv + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(w+1)}{2}}{w} dw + \int_{-\infty}^{\infty} \frac{\cos^2 \frac{n\pi(w+1)}{2}}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{n\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(v-1)}{2}}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(v-1)}{2}}{v^2} dv + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(w+1)}{2}}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2 \frac{n\pi(w+1)}{2}}{w^2} dw \right] & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{(2k-1)\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi v - \frac{\pi v}{2} - k\pi + \frac{\pi}{2} \right)}{v} dv + \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi v - \frac{\pi v}{2} - k\pi + \frac{\pi}{2} \right)}{v^2} dv \right. \\ \quad \left. + \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi w - \frac{\pi w}{2} + k\pi - \frac{\pi}{2} \right)}{w} dw + \int_{-\infty}^{\infty} \frac{\cos^2 \left(k\pi w - \frac{\pi w}{2} + k\pi - \frac{\pi}{2} \right)}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{(2k)\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v^2} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w^2} dw \right] & \text{if } n = 2k \end{cases}
 \end{aligned}$$

Use the facts that $\cos(x \pm \frac{\pi}{2}) = \mp \sin x$ and $\sin(x \pm k\pi) = \sin x \cos k\pi \pm \cos x \sin k\pi = (-1)^k \sin x$.

$$\langle p^2 \rangle = \begin{cases} \frac{(2k-1)\hbar^2}{a^2} \left[- \int_{-\infty}^{\infty} \frac{[-\sin(k\pi v - \frac{\pi v}{2} - k\pi)]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[-\sin(k\pi v - \frac{\pi v}{2} - k\pi)]^2}{v^2} dv \right. \\ \left. + \int_{-\infty}^{\infty} \frac{[\sin(k\pi w - \frac{\pi w}{2} + k\pi)]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[\sin(k\pi w - \frac{\pi w}{2} + k\pi)]^2}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{(2k)\hbar^2}{a^2} \left[- \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - k\pi)}{v^2} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w + k\pi)}{w^2} dw \right] & \text{if } n = 2k \end{cases}$$

$$= \begin{cases} \frac{(2k-1)\hbar^2}{a^2} \left[- \int_{-\infty}^{\infty} \frac{[-(-1)^k \sin(k\pi v - \frac{\pi v}{2})]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[-(-1)^k \sin(k\pi v - \frac{\pi v}{2})]^2}{v^2} dv \right. \\ \left. + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin(k\pi w - \frac{\pi w}{2})]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin(k\pi w - \frac{\pi w}{2})]^2}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{(2k)\hbar^2}{a^2} \left[- \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi v]^2}{v} dv + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi v]^2}{v^2} dv \right. \\ \left. + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi w]^2}{w} dw + \int_{-\infty}^{\infty} \frac{[(-1)^k \sin k\pi w]^2}{w^2} dw \right] & \text{if } n = 2k \end{cases}$$

v and w are just dummy variables, so the integrals can be cancelled and combined.

$$\begin{aligned}
 \langle p^2 \rangle &= \begin{cases} \frac{(2k-1)\hbar^2}{a^2} \left[-\int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - \frac{\pi v}{2})}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi v - \frac{\pi v}{2})}{v^2} dv \right. \\ \qquad \qquad \qquad \left. + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w - \frac{\pi w}{2})}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w - \frac{\pi w}{2})}{w^2} dw \right] & \text{if } n = 2k - 1 \\ \frac{(2k)\hbar^2}{a^2} \left(-\int_{-\infty}^{\infty} \frac{\sin^2 k\pi v}{v} dv + \int_{-\infty}^{\infty} \frac{\sin^2 k\pi v}{v^2} dv + \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w} dw + \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w^2} dw \right) & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{2(2k-1)\hbar^2}{a^2} \int_{-\infty}^{\infty} \frac{\sin^2(k\pi w - \frac{\pi w}{2})}{w^2} dw & \text{if } n = 2k - 1 \\ \frac{2(2k)\hbar^2}{a^2} \int_{-\infty}^{\infty} \frac{\sin^2 k\pi w}{w^2} dw & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{2(2k-1)\hbar^2}{a^2} \left[\left(k - \frac{1}{2}\right) \pi^2 \right] & \text{if } n = 2k - 1 \\ \frac{2(2k)\hbar^2}{a^2} (k\pi^2) & \text{if } n = 2k \end{cases} \\
 &= \begin{cases} \frac{(2k-1)^2 \pi^2 \hbar^2}{a^2} & \text{if } n = 2k - 1 \\ \frac{(2k)^2 \pi^2 \hbar^2}{a^2} & \text{if } n = 2k \end{cases} \\
 &= \frac{n^2 \pi^2 \hbar^2}{a^2}
 \end{aligned}$$

This is the result found in Problem 2.4 using the position-space wave function.