

Problem 4.73

Example 4.4, couched in terms of forces, was a quasi-classical explanation for the Stern–Gerlach effect. Starting from the Hamiltonian for a neutral, spin-1/2 particle traveling through the magnetic field given by Equation 4.169,

$$H = \frac{p^2}{2m} - \gamma \mathbf{B} \cdot \mathbf{S},$$

use the generalized Ehrenfest theorem (Equation 3.73) to show that

$$m \frac{d^2}{dt^2} \langle z \rangle = \gamma \alpha \langle S_z \rangle.$$

Comment: Equation 4.170 is therefore a correct quantum-mechanical statement, with the understanding that the quantities refer to expectation values.

Solution

Equation 3.73 is on page 110 and holds for any observable Q .

$$\frac{d}{dt} \langle Q \rangle = \frac{i}{\hbar} \langle [\hat{H}, \hat{Q}] \rangle + \left\langle \frac{\partial \hat{Q}}{\partial t} \right\rangle \quad (3.73)$$

Recall also the commutator identities shown in Exercise 3.14 on page 108.

$$\begin{aligned} [\hat{A} + \hat{B}, \hat{C}] &= [\hat{A}, \hat{C}] + [\hat{B}, \hat{C}] \\ [\hat{A}\hat{B}, \hat{C}] &= \hat{A} [\hat{B}, \hat{C}] + [\hat{A}, \hat{C}] \hat{B} \end{aligned}$$

Recall also the commutation relations for the components of \mathbf{r} and \mathbf{p} shown in Exercise 4.1 on page 132.

$$\begin{aligned} [r_j, p_k] &= -[p_j, r_k] = i\hbar \delta_{jk} \\ [r_j, r_k] &= [p_j, p_k] = 0 \end{aligned}$$

Start by using Equation 3.73 to find the time derivative of the expectation value of z .

$$\begin{aligned} \frac{d}{dt} \langle z \rangle &= \frac{i}{\hbar} \langle [\hat{H}, z] \rangle + \overbrace{\left\langle \frac{\partial z}{\partial t} \right\rangle}^{=0} \\ &= \frac{i}{\hbar} \left\langle \left[\frac{p^2}{2m} - \gamma \mathbf{B} \cdot \mathbf{S}, z \right] \right\rangle \\ &= \frac{i}{\hbar} \left\langle \left[\frac{p^2}{2m}, z \right] - [\gamma \mathbf{B} \cdot \mathbf{S}, z] \right\rangle \end{aligned}$$

Bring the constants in front. Note that the magnetic field in Equation 4.169 on page 174 is $\mathbf{B}(x, y, z) = -\alpha x \hat{\mathbf{x}} + (B_0 + \alpha z) \hat{\mathbf{z}}$.

$$\begin{aligned}
 \frac{d}{dt} \langle z \rangle &= \frac{i}{\hbar} \left\langle \frac{1}{2m} [p^2, z] - \gamma [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle [p^2, z] \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle [p_x^2 + p_y^2 + p_z^2, z] \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle [p_x^2, z] + [p_y^2, z] + [p_z^2, z] \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle p_x [p_x, z] + [p_x, z] p_x + p_y [p_y, z] + [p_y, z] p_y + p_z [p_z, z] + [p_z, z] p_z \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle p_x(0) + (0)p_x + p_y(0) + (0)p_y + p_z(-i\hbar) + (-i\hbar)p_z \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{i}{2m\hbar} \left\langle -2i\hbar p_z \right\rangle - \frac{i\gamma}{\hbar} \left\langle [\mathbf{B} \cdot \mathbf{S}, z] \right\rangle \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \langle \Psi | [\mathbf{B} \cdot \mathbf{S}, z] | \Psi \rangle \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} \Psi^\dagger [\mathbf{B} \cdot \mathbf{S}, z] \Psi \, d\mathcal{V} \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} \Psi^\dagger [(\mathbf{B} \cdot \mathbf{S})z - z(\mathbf{B} \cdot \mathbf{S})] \Psi \, d\mathcal{V} \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} (\Psi_+ \chi_+ + \Psi_- \chi_-)^\dagger [(\mathbf{B} \cdot \mathbf{S})z - z(\mathbf{B} \cdot \mathbf{S})] (\Psi_+ \chi_+ + \Psi_- \chi_-) \, d\mathcal{V} \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger [(\mathbf{B} \cdot \mathbf{S})z - z(\mathbf{B} \cdot \mathbf{S})] \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} \, d\mathcal{V} \\
 &= \frac{1}{m} \langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \{ [-\alpha x \mathbf{S}_x + (B_0 + \alpha z) \mathbf{S}_z] z - z [-\alpha x \mathbf{S}_x + (B_0 + \alpha z) \mathbf{S}_z] \} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} \, d\mathcal{V}
 \end{aligned}$$

Continue the simplification.

$$\begin{aligned}
 \frac{d}{dt}\langle z \rangle &= \frac{1}{m}\langle p_z \rangle - \frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \frac{\hbar}{2} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} z - z \frac{\hbar}{2} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
 &= \frac{1}{m}\langle p_z \rangle - \frac{i\gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} z - z \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
 &= \frac{1}{m}\langle p_z \rangle - \frac{i\gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} B_0 z + \alpha z^2 & -\alpha x z \\ -\alpha x z & -B_0 z - \alpha z^2 \end{bmatrix} - \begin{bmatrix} B_0 z + \alpha z^2 & -\alpha x z \\ -\alpha x z & -B_0 z - \alpha z^2 \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
 &= \frac{1}{m}\langle p_z \rangle - \frac{i\gamma}{2} \underbrace{\iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V}}_{=0} \\
 &= \frac{1}{m}\langle p_z \rangle
 \end{aligned}$$

Multiply both sides by m , differentiate both sides with respect to t , and use Equation 3.73 again.

$$\begin{aligned}
 m \frac{d}{dt}\langle z \rangle &= \langle p_z \rangle \\
 \frac{d}{dt} \left(m \frac{d}{dt}\langle z \rangle \right) &= \frac{d}{dt}\langle p_z \rangle \\
 m \frac{d^2}{dt^2}\langle z \rangle &= \frac{i}{\hbar} \langle [\hat{H}, p_z] \rangle + \overbrace{\left\langle \frac{\partial p_z}{\partial t} \right\rangle}^{=0} \\
 &= \frac{i}{\hbar} \left\langle \left[\frac{p^2}{2m} - \gamma \mathbf{B} \cdot \mathbf{S}, p_z \right] \right\rangle \\
 &= \frac{i}{\hbar} \left\langle \left[\frac{p^2}{2m}, p_z \right] - [\gamma \mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle \\
 &= \frac{i}{\hbar} \left\langle \frac{1}{2m} [p^2, p_z] - \gamma [\mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle \\
 &= \frac{i}{\hbar} \left\langle \frac{1}{2m} [p_x^2 + p_y^2 + p_z^2, p_z] - \gamma [\mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle
 \end{aligned}$$

Continue the simplification, noting that the magnetic field in Equation 4.169 on page 174 is $\mathbf{B}(x, y, z) = -\alpha x \hat{\mathbf{x}} + (B_0 + \alpha z) \hat{\mathbf{z}}$.

$$\begin{aligned}
m \frac{d^2}{dt^2} \langle z \rangle &= \frac{i}{\hbar} \left\langle \frac{1}{2m} ([p_x^2, p_z] + [p_y^2, p_z] + [p_z^2, p_z]) - \gamma [\mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle \\
&= \frac{i}{\hbar} \left\langle \frac{1}{2m} (p_x [p_x, p_z] + [p_x, p_z] p_x + p_y [p_y, p_z] + [p_y, p_z] p_y + p_z [p_z, p_z] + [p_z, p_z] p_z) - \gamma [\mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle \\
&= \frac{i}{\hbar} \left\langle \frac{1}{2m} [p_x(0) + (0)p_x + p_y(0) + (0)p_y + p_z(0) + (0)p_z] - \gamma [\mathbf{B} \cdot \mathbf{S}, p_z] \right\rangle \\
&= -\frac{i\gamma}{\hbar} \langle [\mathbf{B} \cdot \mathbf{S}, p_z] \rangle \\
&= -\frac{i\gamma}{\hbar} \langle \Psi | [\mathbf{B} \cdot \mathbf{S}, p_z] | \Psi \rangle \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} \Psi^\dagger [\mathbf{B} \cdot \mathbf{S}, p_z] \Psi d\mathcal{V} \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} \Psi^\dagger [(\mathbf{B} \cdot \mathbf{S})p_z - p_z(\mathbf{B} \cdot \mathbf{S})] \Psi d\mathcal{V} \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} (\Psi_+ \chi_+ + \Psi_- \chi_-)^\dagger [(\mathbf{B} \cdot \mathbf{S})p_z - p_z(\mathbf{B} \cdot \mathbf{S})] (\Psi_+ \chi_+ + \Psi_- \chi_-) d\mathcal{V} \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger [(\mathbf{B} \cdot \mathbf{S})p_z - p_z(\mathbf{B} \cdot \mathbf{S})] \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \{ [-\alpha x \mathbf{S}_x + (B_0 + \alpha z) \mathbf{S}_z] p_z - p_z [-\alpha x \mathbf{S}_x + (B_0 + \alpha z) \mathbf{S}_z] \} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
&= -\frac{i\gamma}{\hbar} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \frac{\hbar}{2} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} p_z - p_z \frac{\hbar}{2} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
&= -\frac{i\gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \left(-i\hbar \frac{\partial}{\partial z} \right) - \left(-i\hbar \frac{\partial}{\partial z} \right) \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
&= \frac{\hbar\gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ - \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \frac{\partial}{\partial z} + \frac{\partial}{\partial z} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \right\} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V}
\end{aligned}$$

Therefore,

$$\begin{aligned}
m \frac{d^2 \langle z \rangle}{dt^2} &= \frac{\hbar \gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} -B_0 - \alpha z & \alpha x \\ \alpha x & B_0 + \alpha z \end{bmatrix} \begin{bmatrix} \frac{\partial \Psi_+}{\partial z} \\ \frac{\partial \Psi_-}{\partial z} \end{bmatrix} + \frac{\partial}{\partial z} \begin{bmatrix} B_0 + \alpha z & -\alpha x \\ -\alpha x & -B_0 - \alpha z \end{bmatrix} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} \right\} d\mathcal{V} \\
&= \frac{\hbar \gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} (-B_0 - \alpha z) \frac{\partial \Psi_+}{\partial z} + \alpha x \frac{\partial \Psi_-}{\partial z} \\ \alpha x \frac{\partial \Psi_+}{\partial z} + (B_0 + \alpha z) \frac{\partial \Psi_-}{\partial z} \end{bmatrix} + \frac{\partial}{\partial z} \begin{bmatrix} (B_0 + \alpha z) \Psi_+ - \alpha x \Psi_- \\ -\alpha x \Psi_+ + (-B_0 - \alpha z) \Psi_- \end{bmatrix} \right\} d\mathcal{V} \\
&= \frac{\hbar \gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \left\{ \begin{bmatrix} (-B_0 - \alpha z) \frac{\partial \Psi_+}{\partial z} + \alpha x \frac{\partial \Psi_-}{\partial z} \\ \alpha x \frac{\partial \Psi_+}{\partial z} + (B_0 + \alpha z) \frac{\partial \Psi_-}{\partial z} \end{bmatrix} + \begin{bmatrix} (B_0 + \alpha z) \frac{\partial \Psi_+}{\partial z} + \alpha \Psi_+ - \alpha x \frac{\partial \Psi_-}{\partial z} \\ -\alpha x \frac{\partial \Psi_+}{\partial z} + (-B_0 - \alpha z) \frac{\partial \Psi_-}{\partial z} - \alpha \Psi_- \end{bmatrix} \right\} d\mathcal{V} \\
&= \frac{\hbar \gamma}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \begin{bmatrix} \alpha \Psi_+ \\ -\alpha \Psi_- \end{bmatrix} d\mathcal{V} \\
&= \frac{\hbar \gamma \alpha}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \begin{bmatrix} \Psi_+ \\ -\Psi_- \end{bmatrix} d\mathcal{V} \\
&= \frac{\hbar \gamma \alpha}{2} \iiint_{\text{all space}} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix}^\dagger \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \Psi_+ \\ \Psi_- \end{bmatrix} d\mathcal{V} \\
&= \gamma \alpha \iiint_{\text{all space}} (\Psi_+ \chi_+ + \Psi_- \chi_-)^\dagger S_z (\Psi_+ \chi_+ + \Psi_- \chi_-) d\mathcal{V} \\
&= \gamma \alpha \iiint_{\text{all space}} \Psi^\dagger S_z \Psi d\mathcal{V} \\
&= \gamma \alpha \langle \Psi | S_z | \Psi \rangle \\
&= \gamma \alpha \langle S_z \rangle.
\end{aligned}$$